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Abstract

In this correspondence our aim is to use some tight lower and upper bounds for the differential quaternary phase shift keying transmission bit error rate in order to deduce accurate approximations for the bit error rate by improving the known results in the literature. The computation of our new approximate expressions are significantly simpler than that of the exact expression.

Index Terms

Bit error rate, Marcum Q-function, bounds.

I. INTRODUCTION

Recently, Ferrari and Corazza [\[1\]](#page-6-0) have studied the performance analysis of the differential quaternary phase shift keying (DQPSK) transmission with Gray coding over the additive white Gaussian noise (AWGN) channel. We note that the bit error rate (BER) in the case of Gray coding can be written in terms of the Marcum Q-function and the modified Bessel function of the first kind and zero order. In [\[1\]](#page-6-0) the authors used some bounds for the Marcum Q -function in order to propose some simple, but accurate approximations for the BER. By using some other bounds for the Marcum Q-function (deduced by Wang and Wu [\[3\]](#page-6-1), and Baricz and Sun [\[2\]](#page-6-2)),

The research of Á. Baricz was supported by the János Bolyai Research Scholarship of the Hungarian Academy of Sciences and by the Romanian National Council for Scientific Research in Education CNCSIS-UEFISCSU, project number PN-II-RU-PD 388/2012. The authors are indebted to their friend Yin Sun for his useful comments.

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very recently Sun et al. [\[4\]](#page-6-3) proposed another approximation for BER. In this correspondence we make a contribution to the subject by using some new results on Marcum Q-function: we deduce a new tight upper bound for the BER and combining this with some existing tight bounds we construct some new accurate approximations. These results improve and complement the results of Ferrari and Corazza [\[1\]](#page-6-0), and also of Sun et al. [\[4\]](#page-6-3). In Section [II](#page-1-0) we discuss the bounds for BER of DQPSK, in Section [III](#page-3-0) we present five new approximations for BER and we compare them with the similar ones deduced from inequalities for the Marcum Q-function. Finally, in Section [IV](#page-5-0) the conclusion of the paper is given.

II. BOUNDS FOR BER OF DQPSK

It is known that for the DQPSK transmission with Gray coding over an AWGN channel the BER can expressed as follows [\[5\]](#page-6-4):

BER
$$
\triangleq Q(a, b) - \frac{1}{2}I_0(ab)e^{-\frac{a^2+b^2}{2}}
$$
, (1)

where $a = \sqrt{\gamma(2-\sqrt{2})}$, $b = \sqrt{\gamma(2+\sqrt{2})}$ and γ is the bit signal-to-noise ratio (SNR). Here $Q(a, b)$ stands for the Marcum Q-function, defined by

$$
Q(a,b) \triangleq \int_b^{\infty} x e^{-\frac{x^2 + a^2}{2}} I_0(ax) dx,
$$

where $b \ge 0$, $a > 0$ and I_0 is the modified Bessel function of the first kind and zero order. Note that the Marcum Q-function play an important role in the studies of digital communications over fading channels [\[6\]](#page-6-5).

Now, let us consider the complementary error function, defined by

$$
\operatorname{erfc}(x) \triangleq \frac{2}{\sqrt{\pi}} \int_x^{\infty} e^{-t^2} dt,
$$

and for shortness let us introduce the notation

$$
e(a,b) \triangleq \text{erfc}\left(\frac{b-a}{\sqrt{2}}\right) = \frac{2}{\sqrt{\pi}} \int_{\frac{b-a}{\sqrt{2}}}^{\infty} e^{-t^2} dt.
$$

Recently, Ferrari and Corazza [\[1\]](#page-6-0) by using their bounds from [\[7\]](#page-6-6) (see also [\[8\]](#page-6-7)) for the Marcum Q-function

$$
Q(a,b) \ge \sqrt{\frac{\pi}{2}} \frac{bI_0(ab)}{e^{ab}} e(a,b),\tag{2}
$$

$$
Q(a,b) \le \frac{I_0(ab)}{e^{ab}} \left[e^{-\frac{(b-a)^2}{2}} + a\sqrt{\frac{\pi}{2}}e(a,b) \right],
$$
 (3)

where $b \ge a > 0$, deduced the following bounds for BER of DQPSK with Gray coding:

$$
BER > I_0(ab) \left[\sqrt{\frac{\pi}{2}} \frac{b}{e^{ab}} e(a, b) - \frac{1}{2} e^{-\frac{a^2 + b^2}{2}} \right] \triangleq L_1,
$$

$$
BER < I_0(ab) \left[\sqrt{\frac{\pi}{2}} \frac{a}{e^{ab}} e(a, b) + \frac{1}{2} e^{-\frac{a^2 + b^2}{2}} \right] \triangleq U_1.
$$

It is important to mention here that motivated by the above results of Ferrari and Corazza, very recently Wang and Wu [\[3\]](#page-6-1) proved that

$$
Q(a,b) \ge \sqrt{\frac{\pi}{2}} \frac{bI_0(ab)}{e^{ab} - e^{-ab}} E(a,b),
$$

while Baricz and Sun [\[2\]](#page-6-2) proved that

$$
Q(a,b) \le \frac{I_0(ab)}{e^{ab} + e^{-ab}} \left[e^{-\frac{(b-a)^2}{2}} + e^{-\frac{(b+a)^2}{2}} + a\sqrt{\frac{\pi}{2}} E(a,b) \right],
$$

where in both of the inequalities $b \ge a > 0$ and

$$
E(a, b) \triangleq e(a, b) - e(-a, b) = \frac{2}{\sqrt{\pi}} \int_{\frac{b-a}{\sqrt{2}}}^{\frac{b+a}{\sqrt{2}}} e^{-t^2} dt.
$$

Now, by using the above lower and upper bounds for the Marcum Q-function, Sun et al. [\[4\]](#page-6-3) deduced the following:

$$
BER > I_0(ab) \left[\sqrt{\frac{\pi}{2}} \frac{b}{e^{ab} - e^{-ab}} E(a, b) - \frac{1}{2} e^{-\frac{a^2 + b^2}{2}} \right] \triangleq L_2,
$$

$$
BER < I_0(ab) \left[\sqrt{\frac{\pi}{2}} \frac{a}{e^{ab} + e^{-ab}} E(a, b) + \frac{1}{2} e^{-\frac{a^2 + b^2}{2}} \right] \triangleq U_2.
$$

We note that the lower and upper bounds for the Marcum Q-function of Wang and Wu [\[3\]](#page-6-1) and of Baricz and Sun [\[2\]](#page-6-2), mentioned above, are tighter than [\(2\)](#page-1-1) and [\(3\)](#page-1-2), as it was pointed out in [\[2\]](#page-6-2), [\[3\]](#page-6-1). This in turn implies that the bounds L_2 and U_2 are tighter than the bounds L_1 and U_1 , that is, we have $L_2 > L_1$ and $U_2 < U_1$.

Now, let us consider the following inequality

$$
Q(a,b) \le \frac{I_0(ab)}{e^{ab} + \lambda_0} \left[a\sqrt{\frac{\pi}{2}}e(a,b) + (e^{ab} + \lambda_0)e^{-\frac{a^2 + b^2}{2}} \right],
$$

which holds for all $b \ge a > 0$ and it was deduced by Baricz [\[9\]](#page-6-8). Here $\lambda_0 = e^{\rho_0} (I_0(\rho_0)/I_1(\rho_0) -$ 1) $\simeq 3.03442206626763$, $\rho_0 \simeq 1.54512596391949$ is the unique simple positive root of the equation $(x + 1)I_1(x) = xI_0(x)$, and I_1 stands for the modified Bessel function of order 1. We note that in the above inequality the constant λ_0 is best possible, that is, cannot be replaced by any larger constant. Since the above inequality is an improvement of [\(3\)](#page-1-2), its equivalent form

$$
BER < I_0(ab) \left[\sqrt{\frac{\pi}{2}} \frac{a}{e^{ab} + \lambda_0} e(a, b) + \frac{1}{2} e^{-\frac{a^2 + b^2}{2}} \right] \triangleq U_3
$$

clearly provides a better upper bound than [\(3\)](#page-1-2), that is, we have $U_3 < U_1$. The lower and upper bounds L_1 , L_2 , U_1 , U_2 and U_3 are shown in Fig. [1](#page-7-0) together with the exact BER as functions of the SNR on the interval $(0, 1.5)$. Surprisingly, all of these bounds are extremely tight.

III. APPROXIMATIONS FOR BER OF DQPSK

Observe that there is an interesting formal symmetry between the bounds L_1 and U_1 , as well as between the bounds L_2 and U_2 . Based on the tight bounds L_1 and U_1 , and in view of the formal symmetry, Ferrari and Corazza [\[1\]](#page-6-0) pointed out that it is natural to approximate the BER simply by considering the arithmetic mean of the quantities L_1 and U_1 . In this spirit, they derived the following approximate expression for BER:

$$
\text{BER}_1 \triangleq \mathcal{A} \left(L_1, U_1; 0.5 \right),
$$

i.e.

$$
BER_1 = \sqrt{\frac{\pi}{8}} \frac{a+b}{e^{ab}} I_0(ab)e(a,b),
$$

where $A(x, y; \omega) = \omega x + (1 - \omega)y$ is the weighted arithmetic mean of x and y with weights ω , 1 – $\omega > 0$. By using the same idea, it is natural to propose the following new approximate expressions for BER:

$$
\text{BER}_2 \triangleq \mathcal{A}\left(L_2, U_2; 0.5\right)
$$

and

$$
\text{BER}_3 \triangleq \mathcal{A} \left(L_2, U_3; 0.5 \right),
$$

that is,

$$
BER_2 = \sqrt{\frac{\pi}{8}} \frac{(a+b)e^{ab} - (a-b)e^{-ab}}{e^{2ab} - e^{-2ab}} I_0(ab)E(a,b)
$$

and

$$
BER_3 = \sqrt{\frac{\pi}{8}} I_0(ab) \left[\frac{bE(a,b)}{e^{ab} - e^{-ab}} + \frac{ae(a,b)}{e^{ab} + \lambda_0} \right].
$$

We note that these approximations are better than BER_1 . A similar approximation to that of $BER₂$ was proposed recently by Sun et al. [\[4\]](#page-6-3) in the form

$$
\text{BER}_4 = \frac{e^{-\frac{(b+a)^2}{2}}}{\sqrt{8\pi ab}} + \frac{1}{4}\left(\sqrt{\frac{a}{b}} + \sqrt{\frac{b}{a}}\right)E(a,b),
$$

however, this is better than BER₁ only for $\gamma \in (0.9, 1.4)$. Motivated by the above results, in this correspondence our aim is to propose the following approximations based on better weight functions ω_i , $i \in \{5, 6, 7\}$:

$$
\begin{aligned} \text{BER}_5 &\triangleq \mathcal{A}\left(L_1, U_1; \omega_5\right), \\ \text{BER}_6 &\triangleq \mathcal{A}\left(L_2, U_2; \omega_6\right), \\ \text{BER}_7 &\triangleq \mathcal{A}\left(L_2, U_3; \omega_7\right), \end{aligned}
$$

where

$$
\omega_{5}(\gamma) = \begin{cases}\n0.65 \sqrt[4]{\gamma}, & \text{if } \gamma < 1 \\
0.5 + 1.1 \cdot \frac{e^{\frac{-\pi}{2\sqrt{\gamma}}}}{\gamma^{1.5}} \cdot \sqrt{\frac{1}{2}}, & \text{if } 1 \le \gamma\n\end{cases},
$$
\n
$$
\omega_{6}(\gamma) = \begin{cases}\n e^{-\frac{\gamma^{2}}{2.9}} \cdot 0.25 + 0.5, & \text{if } 0 \le \gamma < 1 \\
\frac{e^{-\frac{1}{2\gamma + 1}}}{(\gamma + 0.5)^{1.5}} \cdot \sqrt{\frac{1}{2\pi}} \cdot 1.15 + 0.5, & \text{if } 1 \le \gamma < 5 \\
\frac{1}{\pi} (1 + \gamma)^{-1} \cdot 0.65 + 0.5, & \text{if } 5 \le \gamma\n\end{cases}
$$
\n
$$
\omega_{7}(\gamma) = \begin{cases}\n(1 - \gamma)^{2} \cdot 0.95, & \text{if } 0 \le \gamma < 1 \\
0.5 - 1.4 \cdot e^{-\gamma^{1.2}} + 0.02, & \text{if } 1 \le \gamma \le 8 \\
\frac{1}{5.2\gamma} + 0.5, & \text{if } 8 < \gamma\n\end{cases}
$$

These approximations can be rewritten as follows

$$
BER_5 = \omega_5(\gamma)L_1 + (1 - \omega_5(\gamma))U_1,
$$

\n
$$
BER_6 = \omega_6(\gamma)L_2 + (1 - \omega_6(\gamma))U_2,
$$

\n
$$
BER_7 = \omega_7(\gamma)L_2 + (1 - \omega_7(\gamma))U_3.
$$

It is important to note here that the computations of approximate BER expressions BER_i , where $i \in \{1, 2, \ldots, 7\}$, completely avoid the computation of the Marcum Q-function. Tables [I](#page-9-0) and [II](#page-10-0) contain an estimate of BER by using the Matlab function marcumq (BER) and the estimates based upon BER_i, where $i \in \{1, \ldots, 7\}$. As we can see the approximations BER_i, where $i \in$

 $\{5, 6, 7\}$, are better than the approximations BER_i, where $i \in \{1, 2, 3, 4\}$. Moreover, the novel approximate expressions BER_i , where $i \in \{5, 6, 7\}$, are consistent with the results of Weinberg [\[10\]](#page-6-9), obtained by using of Monte-Carlo estimators. We note that in our further calculations we considered the values obtained by the marcumq function as the "exact" values for BER.

We give a hint about the construction of the weight functions. By using the values of BER, L_1 , L_2 , U_1 , U_2 and U_3 first we calculated numerically the values of the weight functions. The representation of the initial numerical values for ω_6 can be seen in Fig. [2.](#page-7-1) This looks like the probability density function of the normal distribution but having a heavy tail, so we started to search for expression involved in probability density functions with heavy tails (eg. Student, Lévy, α –stable). By using the same idea the parameters in the expression of the functions ω_i , $i \in \{5, 6, 7\}$, were obtained by numerical experimentations.

Now, in order to estimate the tightness of the new approximate expression and to compare with the tightness of the known approximate expressions, for $i \in \{5, 6, 7\}$ let us consider the relative errors $\varepsilon_i \triangleq (BER_i - BER) / BER$. Table [III](#page-11-0) contains the relative errors ε_i , where $i \in$ $\{5, 6, 7\}$. Finally, in Fig. [3](#page-8-0) and [4](#page-8-1) the exact BER and the approximate expressions BER_i, where $i \in \{1, \ldots, 7\}$, are shown as functions of the SNR on the intervals $(0, 2]$ and $[2, 4]$. We find that our new approximate expressions are quite accurate for the whole region of SNR.

IV. CONCLUSION

In this correspondence we deduced a new tight upper bound for the BER of DQPSK with Gray coding over the AWGN channel by using a new tight bound deduced for the Marcum Q-function. This and the bounds from [\[2\]](#page-6-2) show that approximations of BER for $\gamma = 1$ obtained by using adaptive Simpson quadrature (see [\[10\]](#page-6-9)) and some previous approximations are not in the interval determined by the lower and upper bounds. Although the estimations from [\[1\]](#page-6-0) and [\[4\]](#page-6-3) are the best possible estimations for large values of SNR, by using constants in the convex combinations of the lower and upper bounds, they can be improved by considering weight functions in the calculation of approximations. The advantage of this technique is that we can obtain better approximations also for small values of SNR. Based on numerical experiments we proposed five new approximate BER expressions for the AWGN channel, which use the best known bounds and which are accurate in the whole region of SNR.

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Fig. 1. BER with lower and upper bounds

Fig. 2. The graph of the weight function ω_6

Fig. 3. Approximations for BER on [0, 2]

Fig. 4. Approximations for BER on [2, 4]

γ (dB)	BER	BER ₁	BER ₂	BER ₃
$\mathbf{1}$	1.639×10^{-1}	1.739×10^{-1}	1.731×10^{-1}	1.556×10^{-1}
2	7.161×10^{-2}	7.324×10^{-1}	7.322×10^{-2}	7.007×10^{-2}
3	3.422×10^{-2}	3.458×10^{-2}	3.458×10^{-2}	3.416×10^{-2}
$\overline{4}$	1.701×10^{-2}	1.711×10^{-2}	1.711×10^{-2}	1.706×10^{-2}
5	8.648×10^{-3}	8.683×10^{-3}	8.683×10^{-3}	8.677×10^{-3}
6	4.461×10^{-3}	4.474×10^{-3}	4.4745×10^{-3}	4.473×10^{-3}
τ	2.325×10^{-3}	2.330×10^{-3}	2.3308×10^{-3}	2.33072×10^{-3}
8	1.221×10^{-3}	1.224×10^{-3}	1.22405×10^{-3}	1.22404×10^{-3}
9	6.459×10^{-4}	6.468×10^{-4}	6.46883×10^{-4}	6.46881×10^{-4}
10	3.431×10^{-4}	3.435×10^{-4}	3.43588×10^{-4}	3.43588×10^{-4}
11	1.830×10^{-4}	1.832×10^{-4}	1.83249×10^{-4}	1.83249×10^{-4}
12	9.798×10^{-5}	9.807×10^{-5}	9.80723×10^{-5}	9.80723×10^{-5}

TABLE I BER AND THE ESTIMATES $\mbox{BER}_1,\,\mbox{BER}_2$ AND \mbox{BER}_3

γ (dB)	BER ₄	BER ₅	BER ₆	BER ₇
$\mathbf{1}$	1.484×10^{-1}	1.677×10^{-1}	1.6383×10^{-1}	1.645×10^{-1}
2	6.908×10^{-2}	7.179×10^{-2}	7.162×10^{-2}	7.133×10^{-2}
3	3.348×10^{-2}	3.423×10^{-2}	3.4226×10^{-2}	3.4233×10^{-2}
$\overline{4}$	1.677×10^{-2}	1.7014×10^{-2}	1.7017×10^{-2}	1.7036×10^{-2}
5	8.5931×10^{-3}	8.6500×10^{-3}	8.6500×10^{-3}	8.6593×10^{-3}
6	4.4788×10^{-3}	4.4624×10^{-3}	4.4616×10^{-3}	4.4651×10^{-3}
7	2.3741×10^{-3}	2.3262×10^{-3}	2.3257×10^{-3}	2.3267×10^{-3}
8	1.2841×10^{-3}	1.2222×10^{-3}	1.2219×10^{-3}	1.2218×10^{-3}
9	7.1447×10^{-4}	6.4613×10^{-4}	6.4597×10^{-4}	6.4594×10^{-4}
10	4.1463×10^{-4}	3.4326×10^{-4}	3.4318×10^{-4}	3.4317×10^{-4}
11	2.5591×10^{-4}	1.8311×10^{-4}	1.8307×10^{-4}	1.8306×10^{-4}
12	1.7151×10^{-4}	9.8011×10^{-5}	9.7990×10^{-5}	9.7989×10^{-5}

TABLE II ESTIMATES OF BER BY USING $\text{BER}_4, \text{BER}_5, \text{BER}_6$ AND BER_7

TABLE III

RELATIVE ERRORS ε_i of the approximate expressions BER_i , where $i \in \{5, 6, 7\}$.