

4-3-2-8-7-6

Dan Freed

Joint work (some in progress) with **Constantin Teleman**

Builds on work with **Mike Hopkins** and **Jacob Lurie**

Dedicated to **Graeme**, whose influence is evident throughout

## The definition of conformal field theory

Graeme Segal

I shall propose a definition of 2-dimensional conformal field theory which I believe is equivalent to that used by Freedman et al. The idea arises from joint work with Quillen.

### §1 The definition

The category  $\mathcal{C}$  is defined as follows. There is a sequence of objects  $\{C_n\}_{n \geq 0}$ , where  $C_n$  is the disjoint union of a set of  $n$

Great theorems are awesome; great definitions are transformational.

# The Definition of Quantum Field Theory

**Definition (Segal):** A 2d conformal field theory is a homomorphism

$$F: \text{Bord}_{\langle 1,2 \rangle}^{\text{conf}} \longrightarrow \mathbf{Vect}_{\mathbb{C}}^{\text{top}}$$

**Definition (Atiyah):** An  $n$ -dimensional topological quantum field theory is a homomorphism

$$F: \text{Bord}_{\langle n-1,n \rangle}^{\mathcal{X}(n)} \longrightarrow \mathbf{Vect}$$

$\mathcal{X}(n)$ : an  $n$ -dimensional (topological) tangential structure

**Remark:**  $\text{Bord}_{\langle n-1,n \rangle}^{SO}$  categorifies the classical bordism group  $\Omega_{n-1}^{SO}$   
TQFT categorifies classical bordism invariants, e.g.

$$\text{Sign}: \Omega_{4k}^{SO} \rightarrow \mathbb{Z}$$

# Invertible Field Theories

$\alpha: \text{Bord}_{\langle n-1, n \rangle} \rightarrow \mathbf{Vect}$  is **invertible** if  $\alpha(M)$  is invertible for every  $M$

$\alpha(X) \in \mathbb{C}^\times$  for  $X^n$  closed

$\alpha(Y)$  is a line for  $Y^{n-1}$  closed

**Example:** There is a 3-dimensional TQFT  $\alpha: \text{Bord}_{\langle 2, 3 \rangle}^{\text{framed}} \rightarrow \mathbf{Vect}$  such that  $\alpha(X) = \mu^{\Theta(X)}$ , where  $\mu = e^{2\pi i/24}$  and  $\Theta(X) \in \Omega_3^{\text{framed}} \cong \mathbb{Z}/24\mathbb{Z}$

Replace codomain with  $\mathcal{C}$ , a symmetric monoidal category

$y \in \mathcal{C}$  is invertible if there exist  $y' \in \mathcal{C}$  and  $(y \otimes y' \xrightarrow{\cong} 1_{\mathcal{C}}) \in \mathcal{C}$

**Remark:** Topological invertible theories factor to a map of spectra

$$\begin{array}{ccc}
 \text{Bord}_{\langle n-1, n \rangle}^{\mathcal{X}(n)} & \xrightarrow{\alpha} & \mathcal{C} \\
 \downarrow & & \uparrow \\
 |\text{Bord}_{\langle n-1, n \rangle}^{\mathcal{X}(n)}| & \xrightarrow{\tilde{\alpha}} & \mathcal{C}^\times
 \end{array}$$

## Extended field theories

Early '90s in connection with topological Chern-Simons theory. Also of interest in non-topological theories, e.g. 2d conformal theories.

**Domain:**  $n$ -category  $\text{Bord}_n$  of bordisms  
 $n$  composition laws on  $n$ -manifolds with corners

**Codomain:** Arbitrary  $n$ -category (or  $(\infty, n)$ -category)

**Example:**  $\mathcal{C} = \text{Alg}$  2-category of algebras, bimodules, intertwiners  
 $\text{Alg} \longrightarrow \mathbf{Cat}$  maps  $A \longmapsto {}_A\text{Mod}$   
+(category number) via algebra structure:  $\text{Alg} = E_1(\mathbf{Vect})$   
 $A$  invertible  $\iff A$  central simple

**Example:**  $G$  finite group.  $A = \text{Map}(G, \mathbb{C})$  under convolution.

$$\alpha: \text{Bord}_2 \longrightarrow \text{Alg}$$

with  $\alpha(\text{pt}) = A$ . Then  $\alpha(S^1) = \text{class functions}$ . **N.B.:**  $\alpha$  is not invertible.  
This is an extended version of 2d **Dijkgraaf-Witten** theory.

# The Cobordism Hypothesis

Applies to fully extended *topological* theories

$\text{Bord}_n^{SO}$ :  $(\infty, n)$ -category of oriented bordisms

$\mathcal{C}$ : symmetric monoidal  $(\infty, n)$ -category

$$F: \text{Bord}_n^{SO} \longrightarrow \mathcal{C}$$

**Cobordism hypothesis (Baez-Dolan-Hopkins-Lurie):**

$F$  is determined by  $F(\text{pt}_+)$ . Furthermore, any  $n$ -dualizable,  $SO_n$ -invariant object  $c \in \mathcal{C}$  determines a theory  $F$  with  $F(\text{pt}_+) = c$ .

$n$ -dualizability: data attached to Morse handles exists

$SO_n$ -invariance: extra data on  $c$

**Example:**  $n = 2$ ,  $\mathcal{C} = \text{Alg}_k$  the **Morita** 2-category of algebras

$A \in \mathcal{C}$  is 2-dualizable if it is finite dimensional semisimple

$SO_2$ -invariance data: Frobenius structure (nondegenerate trace)

## Criteria for Invertibility

**Theorem:** Suppose  $\alpha: \text{Bord}_n^{SO} \rightarrow \mathcal{C}$ . Then if either

- 1  $\alpha(S^k)$  is invertible for some  $k \leq n/2$ ; or
- 2  $\alpha(S^n)$  is invertible and  $\alpha(S^p \times S^{n-1-p})$  is invertible for all  $p$ ,

then  $\alpha$  is invertible.

This is a kind of localization theorem for  $\text{Bord}_n^{SO}$ : e.g., 1 says if we invert  $S^k$  then we invert every bordism.

**Example:**  $n = 2, k = 1, \mathcal{C} = \text{Alg}_k$ . If  $A$  is a 2-dualizable (finite dimensional, semisimple) Frobenius algebra, then it defines  $\alpha: \text{Bord}_2^{SO} \rightarrow \text{Alg}_k$  with  $\alpha(S^1)$  equal to the center of  $A$ . So  $\alpha$  is invertible if the center of  $A$  is  $k$ .

# Proof Sketch

First, by the cobordism hypothesis (easy part) it suffices to prove that  $\alpha(\text{pt}_+)$  is invertible; '+' denotes the orientation. We omit ' $\alpha$ ' and simply say ' $\text{pt}_+$  is invertible'.

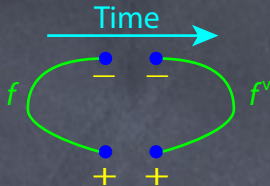
We aim to prove that the 0-manifolds  $\text{pt}_+$  and  $\text{pt}_-$  are inverse:

$$S^0 = \text{pt}_+ \amalg \text{pt}_- = \text{pt}_+ \otimes \text{pt}_- \cong \emptyset^0 = 1$$

with inverse isomorphisms given by

$$f = D^1 : 1 \longrightarrow S^0$$

$$f^\vee = D^1 : S^0 \longrightarrow 1$$



We are reduced to a statement about 1d bordisms: the compositions

$$f^\vee \circ f = S^1 : 1 \longrightarrow 1$$

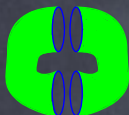
$$f \circ f^\vee : S^0 \longrightarrow S^0$$



must be proved to be identity.



Let's now consider  $n = 2$  where we assume that  $S^1$  is invertible. We apply an easy algebraic lemma which asserts that invertible objects are dualizable and the dualization data is invertible. For  $S^1$  these data are dual cylinders, and so the composition  $S^1 \times S^1$  is also invertible.



**Lemma:** Suppose  $\mathcal{D}$  is a symmetric monoidal category,  $x \in \mathcal{D}$  is invertible, and  $g: 1 \rightarrow x$  and  $h: x \rightarrow 1$  satisfy  $h \circ g = \text{id}_1$ . Then  $g \circ h = \text{id}_x$  and so each of  $g, h$  is an isomorphism.

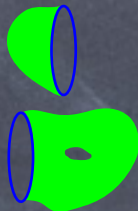
**Proof:**  $x^{-1}$  is a dual of  $x$ ,  $g^\vee = x^{-1}g: x^{-1} \rightarrow 1$ ,  
 $h^\vee = x^{-1}h: 1 \rightarrow x^{-1}$ , so the lemma follows from  $(h \circ g)^\vee = \text{id}_1$ .

Apply the lemma to the 2-morphisms

$$g = D^2: 1 \longrightarrow S^1$$

$$h = S^1 \times S^1 \setminus D^2: S^1 \longrightarrow 1$$

Conclude that  $S^1 \cong 1$  and  $S^2 = g^\vee \circ g$  is invertible.  
 Also,  $g \circ g^\vee = \text{id}_{S^1} \otimes S^2$ , a simple surgery.




Recall that we must prove that the compositions

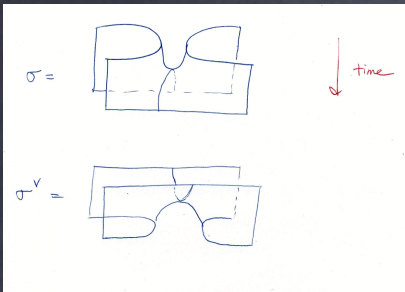
$$f^\vee \circ f = S^1 : 1 \longrightarrow 1$$

$$f \circ f^\vee : S^0 \longrightarrow S^0$$

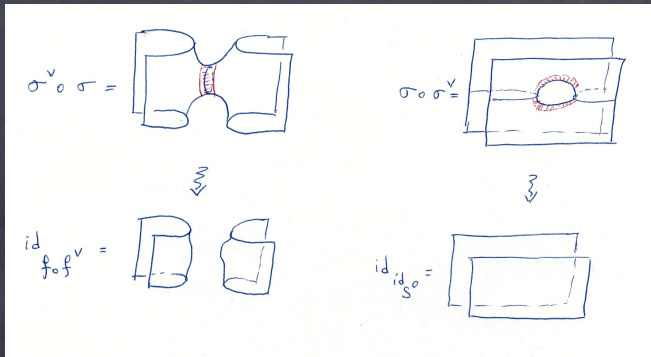
are the identity. We just did the first.



For the second,  $\text{id}_{S^0} =$ 

 and we will show that the saddle  $\sigma : f \circ f^\vee \rightarrow \text{id}_{S^0}$  is an isomorphism with inverse  $\sigma^\vee \otimes S^2$ .



The saddle  $\sigma$  is diffeomorphic to  $D^1 \times D^1$ , which is a manifold with corners. Its dual  $\sigma^\vee$  is the time-reversed bordism.



Inside each composition  $\sigma^v \circ \sigma$  and  $\sigma \circ \sigma^v$  we find a cylinder  $\text{id}_{S^1} = D^1 \times S^1$ , which is  $(S^2)^{-1} \otimes g \circ g^v = (S^2)^{-1} \otimes (S^0 \times D^2)$  by a previous argument. Making the replacement we get the desired isomorphisms to identity maps.

This completes the proof of the theorem in  $n = 2$  dimensions.

In higher dimensions a new ingredient—a dimensional reduction argument—also appears. This uses the Cartesian product on bordisms.

## 4: An Invertible Topological Theory

Let  $A$  be a braided tensor category with braiding  $\beta_{x,y}: x \otimes y \rightarrow y \otimes x$

*Modular tensor category* (quantum group): finiteness and nondegeneracy

**Müger** and others prove nondegeneracy  $\iff$

$$\{x \in A : \beta_{y,x} \circ \beta_{x,y} = \text{id}_{x \otimes y} \text{ for all } y \in A\} = \{\text{multiples of } 1 \in A\} \quad (*)$$

$$= \mathbf{Vect}$$

Braided tensor categories form the objects of a 4-category:

object	category #	
element of $\mathbb{C}$	-1	MTC $A$ : 4-dualizable and $SO_4$ -inv
$\mathbb{C}$ -vector space	0	$\alpha_A: \text{Bord}_4^{SO} \rightarrow \mathbf{Cat}_{\mathbb{C}}^{\beta \otimes}$ ( <b>cob. hyp.</b> )
$\mathbf{Vect}_{\mathbb{C}}$	1	$\alpha_A(\text{pt}_+) = A$
$\mathbf{Cat}_{\mathbb{C}}$	2	$(*) \implies \alpha_A(S^2) = \mathbf{Vect}$ invertible
$\mathbf{Cat}_{\mathbb{C}}^{\otimes} = \mathbf{E}_1(\mathbf{Cat}_{\mathbb{C}})$	3	Thm $\implies \alpha_A$ invertible
$\mathbf{Cat}_{\mathbb{C}}^{\beta \otimes} = \mathbf{E}_2(\mathbf{Cat}_{\mathbb{C}})$	4	<b>Crane-Yetter</b> theory
		$\alpha_A(X^4) = \mu^{\text{Sign}(X)}, \mu = \mu(A) \in \mathbb{C}^{\times}$

**Corollary:** A modular tensor category  $A \in \mathbf{Cat}_{\mathbb{C}}^{\beta \otimes}$  is invertible.

# Relative Quantum Field Theory

**Definition:** Fix an integer  $n \geq 0$  and let  $\alpha$  be an extended  $(n + 1)$ -dimensional quantum field theory. A **quantum field theory  $F$  relative to  $\alpha$**  is a homomorphism

$$F: \mathbf{1} \longrightarrow \tau_{\leq n} \alpha$$

or

$$\tilde{F}: \tau_{\leq n} \alpha \longrightarrow \mathbf{1}$$

**Truncation:** Restrict  $\text{domain}(\alpha)$  along  $\text{Bord}_{\langle n-1, n \rangle} \longrightarrow \text{Bord}_{\langle n-1, n, n+1 \rangle}$

$$F(X^n) \in \alpha(X^n) \quad \text{OR} \quad \begin{aligned} \tilde{F}(X^n; \xi) &\in \mathbb{C}, & \xi &\in \alpha(X^n) \\ \tilde{F}(Y^{n-1}; \mu) &\in \mathbf{Vect}_{\mathbb{C}}^{\text{top}}, & \mu &\in \alpha(Y^{n-1}) \end{aligned}$$

**Representation:**  $\alpha(Y^{n-1}) \longrightarrow \mathbf{Vect}_{\mathbb{C}}^{\text{top}} \quad \alpha(Y^{n-1}) \in \mathbf{Cat}_{\mathbb{C}}^{\text{top}}$

$$\mu \longmapsto \tilde{F}(Y^{n-1}; \mu)$$

## Examples of Relative Theories

**Definition:** If  $\alpha$  is **invertible** we say  $F$  is **anomalous with anomaly  $\alpha$**

$F(X^n) \in \alpha(X^n)$  takes values in a **line** (Lagrangian anomaly)

$F(Y^{n-1}) \in \alpha(Y^{n-1})$  takes values in a **gerbe** (Hamiltonian anomaly)

In  $n = 2$  conformal field theory  $\alpha(X^2) =$  determinant line of Riemann surface  $X$ . Dual picture:  $\tilde{F}(X; \xi) \in \mathbb{C}$  for  $\xi \in \alpha(X)$ .

**Example:**  $n = 1$ ,  $\mathcal{C} = \text{Alg}_k$ ,  $\alpha(\text{pt}) = A = \text{Map}(G, \mathbb{C})$

$V$	finite dimensional representation of $G$
$F(\text{pt}): \mathbf{1} \rightarrow A$	$V$ as a left $A$ -module
$F(S^1) \in \alpha(S^1)$	<b>character</b> of $V$

**Canonical choice:**  $V = A$  with left action (regular representation)

The 2d theory  $\alpha$  only depends on the Morita class of  $A$ , but the relative theory depends on  $A$  itself.

### 3-4: Chern-Simons as a Relative Theory

$A \in \mathbf{Cat}_{\mathbb{C}}^{\beta \otimes}$  modular tensor category  
 $\alpha_A$  4d invertible theory with values in  $\mathbf{Cat}_{\mathbb{C}}^{\beta \otimes} = \mathbf{E}_2(\mathbf{Cat}_{\mathbb{C}})$   
 $F_A$  3d relative theory:  $A$  as left  $A$ -module

$\alpha_A: \mathbf{Bord}_4^{SO} \rightarrow \mathbf{Cat}_{\mathbb{C}}^{\beta \otimes}$  Crane- Yetter theory discussed earlier  
 $F_A: \mathbf{1} \rightarrow \alpha_A$  Chern-Simons theory as a relative oriented theory

**Remarks:**  $F_A$  is defined on the 3d oriented bordism category  $\mathbf{Bord}_3^{SO}$   
 $\alpha_A$  is the usual framing anomaly of Chern-Simons  
Finiteness properties of  $A$  prove that  $F_A$  exists  
Reshetikhin-Turaev recast in terms of cobordism hypothesis  
Walker has a similar picture

**Problem:** To proceed to chiral Wess-Zumino-Witten as a 2-3 theory we would like

$$\alpha_A \xrightarrow{\cong} \mathbf{1}$$

so that  $\chi: \mathbf{1} \xrightarrow{F_A} \alpha_A \xrightarrow{\cong} \mathbf{1}$  is an absolute 3d theory

## Trivializing the Anomaly

Suppose  $X$  is a closed oriented 3-manifold, and we write  $X = \partial W$  for a compact oriented 4-manifold  $W$  with boundary.

The composition

$$\mathbf{1}(X) = \mathbf{1} \xrightarrow{F_A(X)} \alpha_A(X) \xrightarrow{\alpha_A(W)} \alpha_A(\emptyset^3) = \mathbf{1}$$

is multiplication by a number in  $\mathbb{C}$

$W \rightsquigarrow W'$  multiplies this by  $\lambda^{2\pi i c n / 8}$ , where  $n = \text{Sign}(W' \cup_X W)$   
 $c = \text{central charge}$

Signature structure ( $\sigma$ ) makes sense on 1-, 2-, 3-, and 4-manifolds, and every  $(w_1, \sigma)$ -manifold of these dimensions bounds a  $(w_1, \sigma)$ -manifold. Therefore, we recover the Reshetikhin-Turaev 1-2-3-theory, defined on bordisms with a signature structure.



### 3: Chern-Simons as an Absolute Theory

A **tangential structure** using  $p_1$  (**Blanchet-Habegger-Masbaum-Vogel**). It is, in fact, a **stable** tangential structure. If  $M$  is an oriented bordism, a  **$p_1$ -structure** is a lift of a classifying map of  $TM$ :

$$\begin{array}{ccccc} & & BO\langle w_1, p_1 \rangle & & \\ & \nearrow & \downarrow & & \\ M & \xrightarrow{TM} & BO\langle w_1 \rangle & \xrightarrow{p_1} & K(\mathbb{Z}, 4) \end{array}$$

$(w_1, p_1)$ -bordism groups:

$$\Omega_{\{0,1,2,3,4\}}^{(w_1, p_1)} \cong \{\mathbb{Z}, 0, 0, \mathbb{Z}/3\mathbb{Z}, 0\}$$

To define an absolute theory  $\chi$  on  $(w_1, p_1)$ -bordisms we:

- (i) choose a cube root of  $e^{2\pi ic/8}$
- (ii) formally extend the theory to  $\text{pt}_+$  and  $\text{pt}_-$  (not described here)

## 2-3: Chiral WZW as a Relative Theory

$$\alpha_{CA}: \text{Bord}_{\langle 1,2,3 \rangle}^{(w_1, p_1, \text{conf})} \longrightarrow \mathbf{Cat}_{\mathbb{C}} \quad \text{conformal anomaly (invertible)}$$

$$\chi: \text{Bord}_{\langle 1,2,3 \rangle}^{(w_1, p_1)} \longrightarrow \mathbf{Cat}_{\mathbb{C}} \quad \text{Chern-Simons (topological)}$$

Chiral WZW is a relative 2d theory of maps  $\text{Bord}_2^{(w_1, p_1, \text{conf})} \longrightarrow \mathbf{Cat}_{\mathbb{C}}^{\text{top}}$ :

$$F: \chi \longrightarrow \alpha_{CA}$$

This is, of course, a restatement of Segal's **weakly conformal CFT**.

Increasing specificity of modular tensor category  $A$ :

4:  $A$  up to  $E_2$  Morita equivalence

3-4:  $A$  up to  $E_1$  Morita equivalence

2-3:  $A$  with projective representation  $A \longrightarrow \mathbf{Vect}_{\mathbb{C}}^{\text{top}}$   
(projective loop group representations)

## 4-3-2 for Tori

Let  $\Pi$  be a lattice and  $G = T = \Pi \otimes \mathbb{R}/\mathbb{Z}$  the associated torus. A class  $\lambda \in H^4(BT; \mathbb{Z})$  is an even symmetric bilinear form

$$b: \Pi \times \Pi \rightarrow \mathbb{Z},$$

assumed nondegenerate over  $\mathbb{Q}$ .

The kernel of the induced homomorphism  $T \rightarrow T^*$  is a **Pontrjagin self-dual** finite abelian group  $\pi$ , with a quadratic function

$$q: \pi \longrightarrow \mathbb{Q}/\mathbb{Z}$$

$(\pi, q)$  figures in the twisted equivariant  $K$ -theory (**FHT**) as **Verlinde** conjugacy classes, in special modular tensor categories (**Frölich-Karler, Quinn, ...**), and in toral Chern-Simons (**Belov-Moore, FHLT, ...**)

## 4: Finite Path Integral

**Eilenberg-MacLane:**  $(\pi, q)$  determines a homotopy class of maps

$$q: K(\pi, 2) \longrightarrow K(\mathbb{Q}/\mathbb{Z}, 4)$$

$X^4$  closed oriented, get  $q_X: H^2(X; \pi) \longrightarrow \mathbb{Q}/\mathbb{Z}$ . Path integral over  $\pi$ -gerbes reduces to a finite Gauss sum:

$$\begin{aligned}\alpha(X) &= \sum_{\mathcal{G}} \frac{\#H^0(X; F)}{\#H^1(X; F)} e^{2\pi i q_X(\mathcal{G})} \\ &= \frac{\#H^0(X; F)}{\#H^1(X; F)} \sqrt{\#H^2(X; F)} \exp\left[2\pi i (\text{sign } b)(\text{Sign } X)/8\right] \\ &= (\sqrt{\#F})^{\text{Euler } X} \mu^{(\text{sign } b)(\text{Sign } X)} \quad (\mu = e^{2\pi i/8})\end{aligned}$$

Finite path integral gives fully extended 4d theory  $\alpha$  (**F**, **FHLT**)

For example,  $\alpha(Y^3)$  is the vector space of **invariant** functions  $H^2(Y; \pi) \rightarrow \mathbb{C}$ . An automorphism  $a$  of a  $\pi$ -gerbe  $\mathcal{G} \rightarrow Y$  acts on  $\mathbb{C}$  by  $\exp(2\pi i \bar{b}(\mathcal{G}, a))$ , where

$$\bar{b}: H^2(Y; \pi) \times H^1(Y; \pi) \longrightarrow \mathbb{Q}/\mathbb{Z}$$

Since  $\bar{b}$  is nondegenerate, invariant functions vanish away from trivial  $\mathcal{G}$

It follows from ② of theorem that  $\alpha$  is invertible

Finite path integral takes values in “3-algebras”, so will construct relative 3d theory using canonical  $A$  as left  $A$ -module

Chiral WZW for torus is essentially free, but with topological features, so potential construction using Heisenberg groups built from differential cohomology...

## From 4-3-2 to 8-7-6

Arguments from string theory (**Witten, Strominger**) predict the existence of a 6d superconformal field theory, in some ways analogous to chiral WZW.

We call it Theory  $\mathcal{X}$ . (If nothing else, ' $\mathcal{X}$ ' is for 'six'.)

Witten spoke about it at Segal 60.

It is a relative theory. We sketch ideas to approach 8 and 7-8. These are culled from physics, especially as explained to us by Greg Moore.

Construction of 6-7 in general will require new ideas.

## 8-7 and a Bit of 6

The same data  $(\pi, q)$  should give 8 and 7-8.

Starting point is homotopy class of maps

$$q: K(\pi, 4) \longrightarrow K(\mathbb{Q}/\mathbb{Z}, 8)$$

Analogous story (finite path integral, canonical module, ...) should exist, though many points yet to be understood. Perhaps a construction of **non-interacting** 6d Theory  $\mathcal{X}$ , analogous to 2d chiral WZW for tori.

**Frenkel-Kac-Segal** construction: Level 1 2d chiral WZW for **ADE Lie groups** is chiral WZW for maximal torus.

**False analogy:** This is *not* expected for the 6d Theory  $\mathcal{X}$ : there are **interacting** theories based on **ADE Lie algebras**

## 6-7: Theory $\mathcal{X}$

These are the expectations from physics arguments.

**Data:** A triple  $(\mathfrak{g}, b, \Gamma)$  consisting of

- a real Lie algebra  $\mathfrak{g}$  with an invariant inner product  $b$  such that all coroots have square length 2
- a full lattice  $\Gamma \supset \Gamma'$  extending the coroot lattice  $\Gamma'$  such that the inner product is integral and even on  $\Gamma$

Special cases: (1) ADE Lie algebra  $\mathfrak{g}$  (2)  $(\Pi, b)$  with  $\mathfrak{g} = \Pi \otimes \mathbb{R}$

As before, extract  $(\pi, q)$  from this data.

**Expectations:** (i) 7d **topological** QFT  $\alpha_{\mathfrak{g}} = \alpha_{(\mathfrak{g}, b, \Gamma)} = \alpha_{(\pi, q)}$   
(ii) 6d **superconformal** QFT  $\mathcal{X}_{\mathfrak{g}} = \mathcal{X}_{(\mathfrak{g}, b, \Gamma)}$  relative to  $\alpha_{\mathfrak{g}}$



# The New Mother Lode

$$F + *F = 0$$

Theory  $\mathcal{X}$

---

4d

6d

classical PDE

quantum field theory

Riemannian self-duality  
(dimensions 4, 8, ...)

Lorentzian self-duality  
(dimensions 2, 6, 10, ...)

$(G, \lambda)$

$(\mathfrak{g}, b, \Gamma)$

compact Lie group + level

ADE Lie algebra

**BPS, Hitchin, Nahm, ...**

4d  $N = 4$  super Yang-Mills, ...  
geometric representation theory,  
**Khovanov** homology, ...