

Article **FEM Simulation of Fault Reactivation Induced with Hydraulic Fracturing in the Shangluo Region of Sichuan Province**

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Abstract: Hydraulic fracturing operations possess the capacity to induce the reactivation of faults, increasing the risk of fault slip and seismic activity. In this study, a coupled poroelastic model is established to characterize the distribution and movement of fluids within rock formations in the Shangluo region of Sichuan province, China. The effect of hydraulic fracturing projects on the variations of pore pressure and Coulomb effective stress within a high-permeability fault is analyzed. The potential fault-slip mechanism is investigated. The results show that the fault plays different roles for fluid movement, including the barrier, fluid transport channel, and diversion channel, which is related to injection–production schemes. In addition, fluid injection leads to a high probability of fault reactivation. We find that increasing the injection time and fluid injection rate can result in larger slip distances. The injection production scenarios influence the fault-slip mechanism, resulting in a normal fault or reverse fault. However, the arrangement of production wells around the injection can effectively reduce the risk of fault reactivation.

Keywords: fluid–solid coupling; Coulomb failure stress; hydraulic fracturing; seismic simulation; fault slip

1. Introduction

Shale gas, as an unconventional energy source, has gained significant attention. Recently, the occurrence of fault reactivation induced by fluid injection during the stimulation process has been widely concerning for researchers [\[1](#page-15-0)[–3\]](#page-15-1). Previous studies highlighted the influence of both geologic and engineering factors, such as pre-existing fractures and hydraulic fracturing schemes, on fault slip and earthquakes [\[4](#page-15-2)[–6\]](#page-15-3). Fault activity is deemed to be associated with fracturing operations [\[7\]](#page-15-4). Slow and steady slip occurring on major faults in the vicinity of some shale gas wells during fracturing process is observed [\[8](#page-15-5)[,9\]](#page-15-6). Observations show that fault slip leads to casing deformation and the changes in shear stress during fault activity are associated with seismic events [\[10–](#page-15-7)[12\]](#page-16-0). Currently, the Shangluo region in Sichuan has become a mature shale gas development area, where five earthquakes with magnitudes greater than 4.0 ($M_W > 4.0$) occurred during hydraulic fracturing opera-tions [\[13,](#page-16-1)[14\]](#page-16-2). The largest event $(M_W 4.7)$ took place on 28 January 2017. Significant surface deformation due to fault reactivation was observed [\[15\]](#page-16-3). Understanding the mechanisms and control measures of fault reactivation induced by hydraulic fracturing is crucial for seismic monitoring and mitigation during shale gas exploitation.

Fault-slip events are influenced by multiple factors, such as stress state [\[16,](#page-16-4)[17\]](#page-16-5), fluid properties [\[18](#page-16-6)[–20\]](#page-16-7), temperature [\[21](#page-16-8)[,22\]](#page-16-9), and the combined impact of various factors [\[23](#page-16-10)[,24\]](#page-16-11). Cyclical injection and production cause stress perturbations that lead to seismic activities around fault zones [\[22](#page-16-9)[,25\]](#page-16-12). The maximum magnitude of seismic activities occurs during the injection phase of the fracturing fluid. Even small stress changes during the production flowback phase reactivate the faults and generate a large number of seismicity events [\[26](#page-16-13)[,27\]](#page-16-14). Based on laboratory and field experiments on the permeability and frictional properties of

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faults [\[28\]](#page-16-15), it is possible that stress changes during the fracturing process could induce fault reactivation, leading to long-distance fault slips [\[29\]](#page-16-16).

Numerical models for describing fault slip induced by hydraulic fracturing remain a challenge. There are two categories of numerical methods for the simulation of rock fracturing processes: discontinuous methods and continuous methods. The discontinuous methods most commonly used are based on linear elastic fracture mechanics (LEFM), which include cohesive zone models (CZM) and discrete fracture network models (DFN).

DFN is a commonly used discontinuity method. Initially, a semi-analytical numerical model for the calculation of induced stresses along hydraulic fractures was developed [\[30\]](#page-16-17). The influence of fracture initiation location, fluid pressure, fracture dip, and friction factor on fracture slip was discussed. It was found that the region of high-pressure fluid flow has a higher risk of fault reactivation, while the fluid flows at the bottom of the fault could reduce the fault-slip distance. In the DFN model, the effect of fault reactivation is realized by stress perturbation. The DFN model can reflect the characteristics of fluid–solid coupling in the shear process, where crack opening caused by fluid injection can lead to formation slip [\[31\]](#page-16-18). After improving the discrete fracture model, a test was conducted to demonstrate that the method has lower computational cost and more desirable convergence performance than the standard DFM [\[32\]](#page-16-19). Subsequently, to further improve the DFN model [\[33\]](#page-16-20), both static and dynamic induced effects were considered, which could simulate the seismic response in a fractured rock mass. The DFN method approach can represent the injection-induced stress perturbations and pore pressure changes well, but there are still many problems in modelling large deformations and long-distance sliding of faults.

The method of combining continuous and discontinuous method is more widely used for considering excessive displacement. Many researchers have used FLAC3D to simulate hydraulic fracturing under a variety of conditions, including direct injection [\[5,](#page-15-8)[21\]](#page-16-8), staged injection [\[34\]](#page-16-21), and isothermal water injection [\[35\]](#page-16-22). The possibility of fluid injection induced seismicity was discussed. The stress change during fracturing was explained by considering pore pressure diffusion and poroelastic effects. By monitoring the hydraulic behavior in the rock matrix and along the fault, it was found that the fault is more likely to be reactivated when the shear stress is dominant. In addition, the magnitude of removing grouting induced seismic activity is also influenced by the unequal distribution of faultslip time. Rocks with lower modulus show greater slip distance and seismic activity during shear failure. Some researchers focused on the issue of fracture propagation during fluid injection and developed a fully coupled continuous and discontinuous method [\[36,](#page-16-23)[37\]](#page-16-24) to simulate fluid pressure and fracture propagation mode in porous media. The results show that the conversion between fracture propagation modes depends on the injection flow rate.

While stress perturbations within faults during hydraulic fracturing operations have been widely investigated, the mechanisms underlying induced fault slip have not been thoroughly studied. Moreover, predicting the trend of slip motion and understanding the effects of injection remain crucial unresolved issues. In this study, a fully coupled fluid–solid geomechanical model is proposed to simulate the response of a fault, and then the seismic theory is applied to estimate the magnitude of the seismicity. Three important factors affecting the fault slip behavior are discussed: injection time, injection rate, and production scenarios. This research aims to provide a numerical method for assessing induced seismic events and fault slip distance, and to clarify the mechanisms underlying hydraulic fracturing on fault slip events.

2. Methodology

2.1. Coupled Poroelastic Model

The coupled poroelastic model describes how a saturated porous elastic medium responds to changes in stress and pore pressure. According to this model, changes in pore pressure could affect the pore elastic stress of the medium and, conversely, changes in stress could affect the pore pressure [\[38\]](#page-17-0).

The formulation of consolidation in poroelastic media is given as follows. The dis-
placement years u captures the deformation of the solid skeleton, and then the strain placement vector *u* captures the deformation of the solid skeleton, and then the strain tensor ε at a point inside the deformable medium is expressed according to the following assumption [39]: **2. Methodology 2. Methodology**

stress could affect the pore pressure $\overline{}$

sponds to changes in stress and pore pressure. According to this model, changes in pore

$$
\varepsilon = \frac{1}{2} ((\nabla u)^T + (\nabla u))
$$
 (1)

The equilibrium condition of stress field in the deformable medium is provided based on the poroelastic theory: stic theory: ². The equino *2.1. Coupled Poroelastic Model*

$$
\nabla \cdot \sigma + f = 0 \tag{2}
$$

 ϵ is the hody force. The total chrose tensor reads $\tau = \tau + \mathcal{O}$; $(e-e)$ where \mathcal{O} is the 4th order elasticity tensor, and ":" refers to the double-dot tensor product (or double
the 4th order elasticity tensor, and ":" refers to the double-dot tensor product (or double the 4th order elasticity tensor, and α . Teners to the double-dot tensor product (or double
contraction). The elastic strain ε_{el} is the difference between the total strain ε and all inelastic T_{H} and T_{H} and T_{H} and T_{H} are a saturated fluid-filled porous equation σ $\frac{1}{2}$ T_{eff} and σ be an exite stress contribution ϵ_{ℓ} with contributions from initial stresses and viscoelastic stresses. Therefore, Equation (2) is rewritten as where f is the body force. The total stress tensor reads $\sigma = \sigma_{ex} + \mathbb{C} : (\varepsilon - \varepsilon_{inel})$, where $\mathbb C$ is strains ε_{inel} . There may also be an extra stress contribution σ_{ex} with contributions from p and p are deformation of the deformation of the strain φ solid skeleton, and the strain th **2. A. IS the body 2** *2.1. Coupled Poroelastic Model* ial stresses and viscoelastic stresses. Therefore. Equation (2) is rewritten as pressure could affect the pore elastic stress of the medium and, conversely, changes in

$$
\nabla \cdot (\sigma_{ex} + \mathbb{C} : (\varepsilon - \varepsilon_{inel})) + f = 0 \tag{3}
$$

where $\frac{1}{100}$ is the shear modulus, $\frac{1}{100}$ is the velocity field, $\frac{1}{100}$ is the velocity field, $\frac{1}{100}$ The constitutive equations for a saturated fluid-filled porous elastic medium can be expressed as [38,40] \overline{C}

$$
G\,\nabla^2 u + \frac{G}{1-2v}\,\nabla\,\varepsilon_{vol} - \alpha_B\,\nabla\,p = 0\tag{4}
$$

where *G* is the shear modulus, *u* is the velocity field, *v* is Poisson's ratio, ε_{vol} is the volumetric strain, α_B is the Biot coefficient, and ∇p is the applied pressure gradient.

The pore fluid mass conservation equation can be expressed with [41] The constitutive equations for a saturated fluid-filled porous elastic medium can be The pore fluid mass conservation equation can be expressed with [41] **2. Methodology 2. Methodology 2. Methodology 2. Methodology 2.** *Energy Processes**Pennet*

$$
\frac{\partial}{\partial t} \left(\varepsilon_p \rho_f \right) + \nabla \cdot (\rho u_p) = Q_m \tag{5}
$$

 $\frac{1}{2}$ and solid fields, fields, the storage coefficient *S* is defined
the can be calculated from Equation (3). $\frac{1}{\sqrt{2}}$ and $\frac{1}{\sqrt{2}}$ and where ρ_f is fluid density, ε_n is the medium's porosity, and Q_m is the fluid mass source. To as the weighted compressibility of both [42], which can be calculated from Equation (3): where ρ_f is numbered in the first stresses and ζ_m is the number of ζ is rewritten as ζ . where ρ_f is fluid density, ε_p is the medium's porosity, and Q_m is the fluid mass source. To as the weighted compressibility of both [42], which can be calculated from Equation (3): *Energies* **2024**, *17*, x FOR PEER REVIEW 3 of 19 *Energies* **2024**, *17*, x FOR PEER REVIEW 3 of 19 **2. Methodology** *Energies* **2024**, *17*, x FOR PEER REVIEW 3 of 19 **2. Methodology**

$$
S = \varepsilon_p \chi_f + \frac{\partial \varepsilon_p}{\partial p} \tag{6}
$$

$$
\frac{\partial}{\partial t}(\varepsilon_p \rho) = \rho S \frac{\partial p}{\partial t} \tag{7}
$$

Fluid flow in a porous elastic medium is described with Darcy's law. The velocity of the flow in the flow field is calculated with the flow in the flow field is calculated with Fluid flow in a porous elastic medium is described with Darcy's law. The velocity of the flow in the flow field is calculated with $T_{\rm HOM}$ in the flow field is calculated with $T_{\rm HOM}$ the flow in the flow field is calculated with based on the poroelastic theory: stress could affect the pore pressure [38]. \mathbf{S} The formulation of consolidation in the formulation in porous elastic medium is described with Darcy s law. The velocity \mathbf{S} The formulation of consolidation in portfollows of consolidation in portfollows. The velocity of

$$
u_p = -\frac{k}{\mu} \left(\nabla p + \rho_f g \nabla D \right)
$$
 (8)

where u_p is the Darcy's velocity, *k* is the permeability of the medium, μ is the fluid viscosity ∇D is the difference in elevation. ∇D is the difference in elevation. $\mathcal{L}(\mathcal{L})$ ∇D is the difference in elevation.
Therefore, the mass conservation equation is eventually rewritten as ∂ $\mathcal{L} = \mathcal{L} = \mathcal$ iere u_p is the Darcy s velocity, κ is the permeability of the medium, μ is the fluid viscosity, $\mathbf{F} = \mathbf{F} - \mathbf{F} - \mathbf{F} - \mathbf{F}$ is the body force. The total stress tensor reads $\mathbf{F} = \mathbf{F} - \mathbf{F} - \mathbf{F} - \mathbf{F}$ where u_p is the Darcy's velocity, k is the permeability of the medium, μ is the fluid viscosity, α contraction). The elastic strain α is the difference between the difference between the total strain α

 $V D$ is the difference in elevation.
Therefore, the mass conservation equation is eventually rewritten as quation is eventually rewritten as entually rewritten as Therefore, the mass conservation equation is eventually rewritten as Therefore, the mass conservation equation is eventually rewritten as

$$
\rho_f S \frac{\partial p}{\partial t} + \nabla \cdot \rho_f \left[-\frac{k}{\mu} \left(\nabla p + \rho_f g \nabla D \right) \right] = Q_m - \rho_f \alpha_B \frac{\partial}{\partial t} \varepsilon_{vol}
$$
\n(9)

specified by Darcy's law are satisfied. To ensure accuracy, substantial grid ref specified by Darcy's law are statistical. To ensure accuracy, substantial grid reinferment is
required in the vicinity of interfaces with rapidly changing permeability, such as between faults and cap rocks. Failure to meet the accuracy requirements of mesh discretization may
result in anomalous pore pressure oscillations and reduced convergence rates. Assuming that the fluid flow process is a saturated flow process, the flow equations specified by Darcy's law are satisfied. To ensure accuracy, substantial grid refinement is
required in the vicinity of interfaces with rapidly changing permeability such as between result in anomalous pore pressure oscillations and reduced convergence rates. Assuming that the fluid flow process is a saturated flow process, the flow equations $\frac{1}{1}$ is the velocity field, $\frac{1}{1}$ is $\frac{1}{1}$ is $\frac{1}{1}$ is the velocity field, $\frac{1}{1}$ $\frac{1}{1}$ is the velocity field, $\frac{1}{1}$ is $\frac{1}{1}$ is $\frac{1}{1}$ is the velocity field, $\frac{1}{1}$ where $\mathbf{r} = \mathbf{r} \cdot \mathbf{r}$ is the total stress tensor reads tensor reads $\mathbf{r} = \mathbf{r} \cdot \mathbf{r}$ Assuming that the fluid flow process is a saturated flow process, the f required in the vicinity of interfaces with rapidly changing permeability, such as between
Codds and sense the Frihands word the common memission at a found discutive time were where $\mathbf{r} = \mathbf{r} \cdot \mathbf{r}$ is the body force. The total stress tensor reads $\mathbf{r} = \mathbf{r} \cdot \mathbf{r}$ Assuming that the fluid flow process is a saturated flow process, the how equation $\mathbf{u} = \mathbf{u} + \mathbf$ Assuming that the fluid flow process is a saturated flow process, the flow equations $\mathcal{L}(\mathcal{L})$

2.2. Strain–Permeability Model

This study utilizes a function based on porosity and volumetric strain to express the permeability variation in the model. In addition, the approach proposed by Cappa and Rutqvist [\[43\]](#page-17-5) is used to represent the permeability of fracture zones, which links fracture aperture and fracture spacing:

$$
k = \frac{b^3}{12s} \tag{10}
$$

where *k* is the permeability of fault, *b* is the pore diameter, and *s* is the fracture width. The permeability is affected by both tensile and shear fracturing, which adds to the initial fault permeability. Rutqvist et al. (2013) [\[44\]](#page-17-6) used the permeability variation model that took into account the plastic strain along a fault, which is positively correlated with the fault plane:

$$
k = k_0 + k_f = k_0 + A\left(\varepsilon_n - \varepsilon_n^t\right)^3\tag{11}
$$

where k_0 is the initial permeability of fault, *A* is the constant, and ε_n^t is the threshold strain, which represents the strain at the moment when the fault first starts to slip. This study uses ε_n^t = 1 × 10⁻⁴ and *A* = 1 × 10⁻⁵, indicating a potential increase in permeability by three orders of magnitude. This permeability variation is significant compared to the initial fault permeability. As the fractures propagate, pressure diffuses rapidly within the rock mass.

2.3. Coulomb Failure Stress Changes

The Coulomb failure envelope can be estimated based on the rock strength properties, such as the friction angle, cohesion, and tensile strength. Before injection, faults are commonly in either a critical or stable stress condition (Figure [1a](#page-3-0)). However, the significant change in pore pressure caused by the injection of water can induce the reactivation of the fault (Figure [1b](#page-3-0)). When the pressure of fluid injection reaches a certain level, it can cause slip rauit (rigure 1p). When the pressure or fluid injection reaches a certain level, it can cause siip
(shear failure) by exceeding the failure envelope [\[45\]](#page-17-7). Figure [1c](#page-3-0) illustrates the contraction of the Mohr circle, where the change in the vertical principal stress is determined by the variations in applied load pressure and the Biot coefficient, while the change in the horizontal principal stress is related to the rock's Poisson's ratio [21[,46,](#page-17-8)47]. In this situation, the hydraulic fracture does not directly contact the fault, but the rock is elastically deformed during injection. This change in the local stress field near the fault can activate the fault, as shown in Figure 1c, where the Mohr circle approaches the Coulomb failure envelope.

Figure 1. Schematic representation of the Coulomb failure stress (CFS): (a) shear failure diagram before injection; (b) shear failure diagram after injection; (c) Mohr circles with the Coulomb criteria showing failure showing failure envelopes. In (**c**), black circle represents a stable state and red circle represents shear envelopes. In (c) , black circle represents a stable state and red circle represents shear failure. If the pore pressure *P* increases due to fluid injection, the circle will move to the left and reach the failure envelope.

By utilizing the updated stress state, the evolution of Coulomb failure stress (∆*CFS*) during the injection and production processes can be obtained [\[48\]](#page-17-10) with

$$
\Delta CFS = \Delta \tau - \Delta \sigma_n' = (\Delta \tau - \mu_s \Delta \sigma_n) + \mu_s \Delta P_p \tag{12}
$$

$$
\sigma_n = \frac{1}{2}(\sigma_1 + \sigma_3) - \frac{1}{2}(\sigma_1 - \sigma_3)\cos(2\beta)
$$
\n(13)

$$
\sigma'_n = \sigma_n - P_p \tag{14}
$$

$$
\tau = \frac{1}{2}(\sigma_1 - \sigma_3)\sin(2\beta)
$$
\n(15)

where τ is the shear stress, σ_n is the total normal stress, μ_s is the coefficient of static friction, and P_p is the pore pressure. The Coulomb failure criterion relates the frictional strength of the fault to two factors: (1) the Terzaghi effective stress $(\sigma_n - P_p)$ and (2) the coefficient of static friction, μ_s , which is commonly assumed to range between 0.6 and 0.8. It is assumed that the frictional properties along the fault plane are homogeneous and the representative value of the static friction coefficient, *µs* , is taken as 0.6.

3. Numerical Model Table 1. Material properties used in the models.

3.1. Geologic Model and Boundary Conditions

The study area is located in the Shangluo region of Sichuan province, China. A finite element model was established to simulate the injection process with COMSOL Multiphysics 6.0. The model has a length of 24,000 m and a height of 12,000 m, as shown
Provided by the contract of th in Figure [2.](#page-4-0) Three different distances between the injection well and the fault were set,
as shown in Figure 2.0 (a) d = 150 m (b) d = 450 m and (a) d = 750 m. To investigate the as shown in Figure [3a](#page-5-0): (a) $d = 150$ m, (b) $d = 450$ m, and (c) $d = 750$ m. To investigate the influence of injection/production schemes on fault slip, three different configurations were set in the simulations: (1) case 1 with two injection wells; (2) case 2 with one injection well located on the hanging wall of the fault and one production well located on the footwall of the fault; and (3) case 3 with one injection well located on the footwall of the fault and one production well located on the hanging wall of the fault (Figure [3b](#page-5-0)). Note that the spacing
D hetween two wells in each case is fixed at 1500 m *D* between two wells in each case is fixed at 1500 m.

Figure 2. Two-dimensional geological model of the study area. T₁, P, Z, and Pt₂₋₃ represent limestone formation, and S represents shale reservoir.

The sedimentary sequence of the study area consists of a dense, low-permeability limestone caprock, a shale reservoir and a fault. The material properties are presented in Table [1](#page-5-1) [\[49\]](#page-17-11). The hydraulic fracturing target is the Longmaxi shale formation, which is buried at a depth of approximately 2.7 to 3 km. The reservoir has a thickness of approximately 200 m and is restricted by adjacent layers with very low porosity and permeability. The fault, which has a dip angle of 80° , is 40 m thick. It is commonly characterized as comprising a fault core and the surrounding damage zone [\[50](#page-17-12)[–52\]](#page-17-13).

3.2. Boundary Condition and Parameter Setting **Table 1.** Material properties used in the models.

3.2. Boundary Condition and Parameter Setting

The solid mechanics equations are subject to roller support boundary conditions on both the left and right sides of the model. The bottom boundary is fixed, while the top boundary is free. Regarding the fluid flow equations, a non-flow boundary condition is set at the bottom of the model, where the normal component of velocity is zero. To ensure a uniform pressure distribution along the fault, it is important to have a sufficiently wide domain width. The injection and production processes are simulated using well boundary conditions, with injection and production rates of $15 \text{ m}^3/h$ and a total simulation time of 50 h. A finite element mesh is generated using the free triangular grid method.

To replicate the in situ stress conditions in the study area, the model is initialized with the following steps:

- (1) Effective stress initialization: The initial stress state input for the solid mechanics calculations is pore pressure, which represents pre-existing stress conditions.
- (2) To prevent the calculation of induced slip from being affected by settlement effects, it is necessary to apply the stress state induced by gravity settlement to the undeformed model. This can be achieved by performing two iterations of steady-state calculations, using the results from the first iteration as the starting point for the second iteration. This iterative process helps ensure convergence of the calculation and eliminates the influence of settlement displacement on simulation results.

By following these initialization steps, the model aims to capture realistic stress conditions and reduce the impact of settlement displacement on the calculation of induced slip.

inates the influence of settlement displacement on simulation results.

3.3. Verification 3.3. Verification

To verify the model, we compare the numerical and analytical solutions of pore pressure distribution after injection using Darcy's law and Biot's theory. The computational model is depicted in Figure [4,](#page-6-0) using the same boundary conditions and material prop-model is depicted in Figure 4, using the same boundary conditions and material propererties as the geological model. Figure [5a](#page-6-1) shows the comparison between the numerical ties as the geological model. Figure 5a shows the comparison between the numerical and and analytical solutions for the maximum pore pressure. The numerical solution yields a value of 7.07 MPa, while the analytical solution yields a value of 6.71 MPa. Figure [5b](#page-6-1) presents the distribution of pore pressure at the boundary of the model obtained with the analytic and numerical solutions. The numerical model considers the enhanced permeability caused by the injection, allowing fluid pressure to diffuse more smoothly than the analytical solution. Both approaches exhibit excellent agreement, validating the reliability of the numerical simulation. numerical simulation. To verify the model, we compare the numerical and analytical solutions of pore pressure distribution after injection using Darcy's law and Biot's theory. The computational

Figure 4. Validation model. **Figure 4.** Validation model.

Figure 5. Results of analytical and numerical solutions. (a) Variation of pore pressure with time at the monitoring point; (b) pore pressure distribution along the monitoring line after 600 s of fluid injection.

4. Results

4.1. Effect of Injection on Fluid Distribution

4.1.1. Pore Pressure Distribution

Figure [6](#page-7-0) shows the pore pressure distribution in the model. High pore pressure initially appears around the injection well. Once the fluid encounters a fault, it rapidly flows along the fault, while a small amount of the fluid diffuses into the surrounding cap rock. When the fluid reaches the ground (see Figure [6c](#page-7-0)), the pore pressure rapidly diffuses along the fault to the deeper layer of the formation. At the end of the simulation, the fault is filled with fluid and the pore pressure remains stable. The fluid is stored in the fault, creating a high-pressure fluid state along it.

Figure 6. Distribution of pore pressure during fluid injection: (a) $t = 10$ h, (b) $t = 25$ h, (c) $t = 40$ h, and (**d**) $t = 50$ h.

4.1.2. Evolution of Pore Pressure along the Fault 4.1.2. Evolution of Pore Pressure along the Fault

Figur[e 7](#page-8-0) shows the pore pressure distribution along the fault. When $d = 150$ m, pore pressure diffusion along the fault is relatively slow during the pre-injection period $(t = 0-9.5 \text{ h})$, as shown in Figure 7a. [As](#page-8-0) the pore pressure within the fault reaches its peak at $t = 12$ h, the fluid rapidly transports along the fault [$43,53$ $43,53$]. It can be seen that a high pore pressure appears about 3000 m below the surface at *t* = 30 h. At the end of the simulation, pressure appears about 3000 m below the surface at *t* = 30 h. At the end of the simulation, the high-pressure fluid fills the entire fault and then the pore pressure inside the fault remains relatively stable, always slightly lower than the peak pressure inside the fault. When the injection well is far away from the fault $(d = 750 \text{ m})$, the distribution characteristics of the pore pressure show similar features. In the pre-injection period, the fluid flow occurs mainly within the reservoir; thus, the pore pressure within the fault is significantly lower than that in the case of $d = 150$ m.

Figure 7. Variation of pore pressure along the fault with time: (a) $d = 150$ m; (b) $d = 750$ m.

Figure 8 s[ho](#page-8-1)ws the pore pressure distribution along the fault at $t = 50$ h for different injection and production scenarios. When the injection well and the production well is set in the simulation, the pore pressure in the fault significantly reduces, and, in some cases, negative pore pressure even occurs. As shown in Figure 9, the dist[anc](#page-9-0)e between the production well and the fault significantly affects fluid distribution characteristics inside the fault. In case 2, where the production well is located far away from the fault, the fault acts as a conduit due to its high permeability, and fluid within the fault flows rapidly towards the production well. As a result, the pore pressure within the fault is significantly lower than that in case 1. However, if the production well is close to the fault (case 3), the presence of the production well prevents the increase in pore pressure within the fault. The production well, being close to the fault, first establishes hydraulic connectivity. Then, fluid in the surrounding rock formations is rapidly drawn in by the $\frac{1}{2}$ well. In this situation, the fault acts as a barrier for fluid flow, resulting in less noticeable
the fault in the fault. pressure changes within the fault.

Figure 8. Pore pressure distribution along the fault at $t = 50$ h for different scenarios.

Figure 9. Pore pressure distribution for different injection–production scenarios when $d = 150$ m: case 1, (**b**) case 2 and (**c**) case 3. (**a**) case 1, (**b**) case 2 and (**c**) case 3.

4.2. Coulomb Failure Stress (∆CFS) 4.2. Coulomb Failure Stress (∆*CFS)*

Injection activities can significantly affect fault stability by generating high pore pres-Injection activities can significantly affect fault stability by generating high pore pressures along the fault. Previous research suggests that the fault could be unstable and activated when the Coulomb failure stress exceeds $0.05-0.2$ MPa $[2,25]$ $[2,25]$ [. In](#page-16-12) the simulations, two different injection locations and three different injection schemes are considered (Figure [10\)](#page-10-0).

When *d* = 150 m, the pore pressure within the fault is much higher for case 1 than When *d* = 150 m, the pore pressure within the fault is much higher for case 1 than other cases [\(Fig](#page-10-0)ure 10a). In case 1, high Coulomb failure stress occurs within the fault, mostly larger than 1 MPa, indicating that the fault is in a relatively active state. In case 2, where the production well is located far away from the fault, there is a significant reduction in the tion in the Coulomb failure stress, which is reduced by approximately 30% compared to Coulomb failure stress, which is reduced by approximately 30% compared to case 1. The activated region in case 2 is located at the base of the fault. In case 3, where the production well is at the proximal end of the fault, most of the area on the fault is close to the activation state, which is prone to induce slip.

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Figure 10. Variation of Coulomb failure stresses along the fault at $t = 50$ h: (a) $d = 150$ m; (b) $d = 750$ m.

In case 1, the value of ΔCFS is low and the possibility of fault slip is small. As shown in Figure 10b, for cases 2 and 3, it is only at the tip of the fluid flow that there is a higher risk of fault activation. The fluid-filled region returns to a relatively stable stress state. The injection and production scheme helps to reduce the possibility of fault activation. When $d = 750$ m, the activation region of the fault is much smaller than that for $d = 150$ m.

Figure 11 illustrates the distri[but](#page-10-1)ion of Coulomb failure stresses at different stages of fluid injection for $d = 150$ m. The plot shows that the Coulomb failure stress changes primarily along the wellbore, the fracture zone, and the fault. Fracturing creates a failure zone that facilitates fluid redistribution. The high-pressure fluid in the zone continues to flow, redistributing pore pressure from high-pressure to low-pressure areas. High-pressure injection causes deformation of the rock, maintaining relatively high pore pressure within the reservoir and the fault. The pore pressure stored in the fault also causes the fault to remain at a high ∆*CFS* level. fault to remain at a high Δ*CFS* level.

Figure 11. Coulomb failure stress distribution at different time for $d = 150$ m: (a) $t = 10$ h, (b) $t = 20$ h, (**c**) $t = 30$ h, and (**d**) $t = 40$ h.

4.3. Induced Seismic Events

Fault reactivation caused by hydraulic fracturing is one potential factor contributing to seismic activity. When the Coulomb stress on both sides of a fault reaches a critical level, seismic activity could be induced, such as in Harrison County and Poland Township in Ohio [38,54], Alberta [55], and Pohang [5]. Since shale gas development began in 2012, [the](#page-17-15) number of seismic events in Shangluo City has increased significantly. Several earthquakes exceeded the magnitude of 3. The sources of induced seismicity are usually located near the injection position, which is typically around 2 to 3 km. Injection-induced seismicity has certain characteristics compared to natural earthquakes, such as higher source intensity and shallower source depth. Slip nodes with a radius of R range can be calculated as α

4.3.1. Calculation of Seismic Slip Events the same event; otherwise, they are considered to be separate events.

Seismic events can be estimated based on the shear displacement and material proper-Seismic events can be estimated based on the shear displacement and material proper-
ties. As shown in Figure [12,](#page-11-0) the vertices of the mesh triangles serve as the nodes at which events occur. Each node has an associated area of influence, which is defined as one third of the total area of all surrounding triangles. The seismic moment can be calculated with

$$
M_0 = GAd \tag{16}
$$

where *G* is the shear modulus, *A* is the total area of sliding nodes in an earthquake event, and *d* is the average value of plastic slip of all slip nodes. where G is the shear modulus, A is the total area of silding hodes in an earthquake event, α

Figure 12. Seismic event estimation method. **Figure 12.** Seismic event estimation method.

Equation (16) only considers shear failure and ignores seismic events caused by tensile failure. It may overestimate the proportion of large seismic events in the model, as some slip behavior does not result in seismic activity. Therefore, only the nodes that experience seismic slip should be counted as seismic events. Each node is treated as an individual event with a radius of *R*. Slip nodes within this range can be calculated as part of the same event; otherwise, they are considered to be separate events.

The moment magnitude can be calculated with [\[44\]](#page-17-6)

$$
M_w = \frac{2}{3} \log_{10} M_0 - 9.1 \tag{17}
$$

Injection-induced fault activity can be classified into three categories on the basis of slip velocity: creep and slow slip, aseismic slip, and seismic slip [\[3\]](#page-15-1). In this study, slip nodes are considered to be seismic slip only if they satisfy two conditions: (1) the Coulomb failure stress is greater than 0.2 MP, and (2) the slip velocity is greater than 0.2 mm/s. The ∆*CFS* threshold guarantees that rock failure occurs within the defined area, while the critical velocity threshold prevents slow creep and aseismic slip [\[27,](#page-16-14)[56\]](#page-17-17).

4.3.2. Distribution of Seismic Slip Events

Figure [13](#page-12-0) presents simulated seismic events caused by fluid injections. The results show a similarity between the distribution of seismic events and that of pore pressure. During the initial stage of injection, the magnitude of seismic events is relatively small (mostly smaller than 2). Typically, micro-seismic events are first observed near the injection well after 4–6 h of injection. However, there is a significant increase in seismic
exercise with a set of the activity of the activity of the T₁ activity after a period of fluid injection as the fluid enters the fault. The majority of these events are earthquakes with $Mw = 2-3$, but some can even reach $Mw > 4$. Seismic activity First occurs near the injection well and then spreads to both sides of the fault. Seismic events occur at the interface between the reservoir and the caprock due to differences in material properties. The reservoir rock is subject to damage due to high-pressure fluid. In contrast, the hydraulic transformation of caprock with low permeability and good rock properties is negligible, resulting in a significant velocity difference of deformation at the reservoir/caprock interface. As a result, seismic events are mostly concentrated at the interface.

Figure 13. Seismic events. (**a**) Actual seismic events in Shangluo, Sichuan, since 1 January 2015, **Figure 13.** Seismic events. (**a**) Actual seismic events in Shangluo, Sichuan, since 1 January 2015, where seismic symbols are colored by date and scaled by magnitude [57]; (**b**) the distribution of where seismic symbols are colored by date and scaled by magnitude [\[57\]](#page-17-18); (**b**) the distribution of associated seismic events after 20 h of fluid injection; (c) the distribution of seismic events after 50 h of fluid injection. of fluid injection.

5. Discussion 5. Discussion

5.1. Effects of Injection on Formation Deformation

5.1. Effects of Injection on Formation Deformation When the fracturing fluid is injected into the target reservoir, it will inevitably produce pore pressure accumulation, effective stress change, and reservoir deformation [58,59]. To prevent th[e](#page-17-19) deformation of the reservoir from being transmitted to the surface and causing damage to ground facilities, it is necessary to consider the ground uplift caused by the injection of fracturing fluid. Under the influence of high-pressure fluid injection, the rock mass above the reservoir deforms, causing an upward move-
interaction, the rock mass above the reservoir deforms, causing an upward move-ment [\[60\]](#page-17-21). Figure [14](#page-13-0) shows that displacement mainly concentrates near the injection well
and the high annual illing facts Figure 14 shows that displacement mainly concentrates near the injection well and the injection well and the injection well and the injection were near the injection were near the injection were near the injection were nea and the high-permeability fault.

Figure 14. Distribution of displacement at different time: (a) $t = 10$ h, (b) $t = 25$ h, (c) $t = 40$ h, and **)** $t = 50$ **h.**

5.2. Fault Slip

5.2.1. Effect of Injection Time $\frac{1}{2}$ shows the variation in fault slip with time. At the beginning, slipped is $\frac{1}{2}$

Figure 15 shows the variation in fault slip [with](#page-13-1) time. At the beginning, slippage is small. At $t = 40$ h, a significant upward shift occurs. At the end of the simulation, the fault is fully filled with fluid and there is relatively stable slippage between the upper and lower walls of the fault, indicating a trend of reverse fault slippage. With a maximum distance of slippage of about 0.14 m, the slip near the surface is greater than that in the deep rock layer.

Figure 15. Cumulative slippage along the fault at different time. **Figure 15.** Cumulative slippage along the fault at different time.

5.2.2. Effect of Injection Rate

Injection rate is a crucial factor that affects the efficiency of hydraulic fracturing and directly impacts the extent of reservoir stimulation. Higher injection pressure often results in better reservoir modification. However, it also increases the probability of induced fault slip. Figure 16 shows that the effect of injection rate on the fault is not significant when the injection rate is less than 7.5 m³/h. As the injection rate increases, the affected area within the fault expands noticeably. Due to the injection well's proximity to the hanging wall of the fault, the hanging wall experiences larger relative displacements compared to the footwall. When the injection rate is 22.5 m³/h, the slippage increases by approximately 25%. Therefore, an appropriate reduction in the injection rate could minimize the occurrence of fault slip without significantly compromising the effectiveness of reservoir modification.

Figure 16. Cumulative slippage along the fault for different injection rates.

5.2.3. Effect of Injection–Production Schemes For cases 2 and 3, where one injection well and one production well are set, slip dis-

Figure 17a shows the impact of injection–production schemes on slip mechanisms. For the cases with $d = 150$ m and 450 m, where the injection well is closer to the fault, a significant velocity component on the hanging wall results in a reverse fault-slip state. When the distance of the two injection wells to the fault is the same $(d = 750 \text{ m})$, the injection on the footwall is dominant, facilitating fluid flow towards the lower portion of the fault. The hanging wall rock mass undergoes downward sliding, resulting in a negative slip difference, which indicates a normal fault behavior.

Figure 17. Slippage along the fault for (a) different spacing between the wells and the fault, and (b) different production scenarios. different production scenarios.

6. Conclusions (1) Case 1 with two injection wells; (2) case 2 with one injection well located on the hanging wall of the fault and one production well located on the footwall of the fault; and $f(3)$ case 3 with one injection well located on the footwall of the fault and one production well located on the hanging wall of the fault.

For cases 2 and 3, where one injection well and one production well are set, slip distance of −0.38 m and −0.31 m occurs, respectively (Figure [17b](#page-14-1)). If the production $\frac{d}{dt}$ wen is far front the fault (case 2), the fault is in a conduit state and flost of the fluid migrates through the fault to the production well. This weakens the movement trend of the magnates unough the fault to the production well. This wells in the movement tiend of the hault (case 3), the ranging wan of the fault. Then the production well is focated close to the fault (close *b*), the surrounding rock mass contracts towards the well, resulting in a relatively slow movement rate. In summary, due to the presence of the production well, the hanging wall motion $\frac{1}{2}$ The migration of the migration of the migration is component is reduced reculting in a normal fault clin mechanism component is reduced, resulting in a normal fault-slip mechanism. well is far from the fault (case 2), the fault is in a conduit state and most of the fluid

6. Conclusions

This study presents a fully coupled numerical investigation of seismic events and fault-slip mechanisms in the Shangluo region of Sichuan province, China. The distribution of fluid pressure and Coulomb failure stress was simulated using FEM. The findings are listed as follows:

- (1) High-permeability faults display three distinct behaviors under different production schemes: barrier, fluid transport channel, and conduit channel. The faults act as conduits in the absence of producing wells. Where the production well is located far away from the fault, the fault acts as a conduit channel. If the production well is close to the fault, the fault acts as a barrier for fluid flow.
- (2) The migration of high-pressure fluid in the formation is closely related to the degree of rock fracturing. This leads to most of the area on the fault close to the activation state. As the fluid distribution within the fault rock mass tends to be stable, the fault can return to a relatively stable stress state.
- (3) The results show that the displacement mainly occurs near the injection well and the fault. The production time and injection rate affect the distance of fault slip, and the fault slip near the surface is greater than in other places. The injection production scenarios could influence the fault-slip mechanism, resulting in a normal fault or reverse fault.

These results could help to understand seismic events and fault behavior in the Shangluo region and provide valuable insights for reservoir stimulations and seismic hazard assessments in similar geological settings.

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