Nuclear Physics B206 (1982) 496–497

Errata

A.G. Liparteliani, V.A. Monich, Yu.P. Nikitin and G.G. Volkov, Neutral mesons with heavy quarks and mixing angles in the six-quark model, Nucl. Phys. B195 (1982) 425.

In eqs. (13)–(15) a multiple 4 has been omitted.

In eqs. (29), (30), (36), (37) $[(\frac{5}{2}-3\ln(y_3/z_2))\frac{5}{12}+(\frac{5}{12}+3\ln(y_3/z_3))]$ must be replaced by $\{\frac{1}{4}[\ln(y_3/z_3)-\frac{9}{2}]\}$.

In eqs. (26)–(28) and (33)–(35) the coefficients $\frac{11}{9}$ should read $\frac{1}{3}$.

R.. D'Auria and P. Fré, Geometric supergravity in $D=11$ and its hidden supergroup, Nucl. Phys. B201 (1982) 101.

In eq. (5.10) the $\frac{1}{8}$ should be replaced by $\frac{1}{8}i$.

In eq. (5.11) -3 should read -6 and $\frac{3}{8}$ should read $\frac{3}{4}$.

In table 3 for the part concerning the on-shell solution for the curvatures the $\frac{1}{8}$ in the expression for ρ should again be replaced by $\frac{1}{8}i$ and in the expression for R^{ab}, $-\frac{7}{9}$ should be replaced by $+1$ and $+\frac{55}{216}$ by $+\frac{1}{24}$. Also in table 3, in propagation equation (iii) the same correction should be made as in eq. (5.11).

M.T. Grisaru and W. Siegel, Supergraphity (II). Manifestly covariant rules and higher-loop finiteness, Nucl. Phys. B201 (1982) 292.

The "doubling" trick of sect. 4 cannot be applied covariantly in the case where the scalar multiplet is a complex representation of the Yang–Mills group. (However, it can be applied as described to supergravity, and to real representations of the Yang–Mills group.) This is due to the fact that $\bar{\nabla}^2\bar{\eta}$ is then not in the same representation as η, so the operator $\boldsymbol{O}$ is not representation-preserving. As a result, one must use rules *at one loop* which are not expressed manifestly in terms of Γ_A.

Explicitly, in terms of fields $\hat{\eta}(\hat{\bar{\eta}})$ and sources $\bar{J}(\hat{\bar{J}})$ which are chiral (antichiral) with respect to $\bar{D}_{\dot{\alpha}}(D_\alpha)$, we have the following equations of motion in the presence of external super-Yang–Mills:

$$\hat{\boldsymbol{O}}\begin{pmatrix}\hat{\eta}\\ \hat{\bar{\eta}}\end{pmatrix}+\begin{pmatrix}\hat{J}\\ \hat{\bar{J}}\end{pmatrix}=0\,,\qquad \hat{\boldsymbol{O}}=\begin{pmatrix}0 & \bar{D}^2\,\mathrm{e}^{V^*}\\ D^2\,\mathrm{e}^{V} & 0\end{pmatrix}.$$

If we had used covariantly chiral fields, we would have

$$\eta=\mathrm{e}^{\bar{W}}\hat{\eta}\,,$$