

## SOLITARY WAVE SOLUTIONS OF THE NAVIER-STOKES EQUATIONS BY HE'S VARIATIONAL METHOD

by

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*Existence of variational principles for Navier-Stokes equations has been discussing for hundreds of years, but it has not yet been solved. In this study, a new perspective is proposed, which uses a traveling wave transform, so that a variational formulation can be established. Furthermore, the solitary wave solutions are solved by He's variational method.*

**Key word:** Navier-Stokes millennium-prize problem, traveling wave transform, solitary wave solutions, He's variational methods, He-Weierstrass function, variational principle

### Introduction

Any a motion should follow a nature law, the most famous one is the Hamilton principle [1-3], which is a minimum variational principle. Navier-Stokes equations describe the motion of a fluid, which can be expressed:

$$\frac{\partial \xi}{\partial s} + (\xi \cdot \nabla) \xi = -\frac{1}{\rho} \nabla p + \frac{\mu}{\rho} \nabla^2 \xi \quad (1)$$
$$\nabla \xi = 0$$

where  $p$  is the pressure,  $\rho$  – the density,  $\mu$  – the kinematic viscosity, and  $s$  – the time. In 3-D space, velocity vector is  $\xi = \xi(\xi_1, \xi_2, \xi_3)$ , the components of Navier-Stokes equations in  $i$ -,  $j$ -, and  $k$ -directions are given by the following equations:

$$\frac{\partial \xi_1}{\partial s} + \xi_1 \frac{\partial \xi_1}{\partial i} + \xi_2 \frac{\partial \xi_1}{\partial j} + \xi_3 \frac{\partial \xi_1}{\partial k} = \frac{\mu}{\rho} \left( \frac{\partial^2 \xi_1}{\partial i^2} + \frac{\partial^2 \xi_1}{\partial j^2} + \frac{\partial^2 \xi_1}{\partial k^2} \right) - \frac{1}{\rho} \frac{\partial p}{\partial i}$$
$$\frac{\partial \xi_2}{\partial s} + \xi_1 \frac{\partial \xi_2}{\partial i} + \xi_2 \frac{\partial \xi_2}{\partial j} + \xi_3 \frac{\partial \xi_2}{\partial k} = \frac{\mu}{\rho} \left( \frac{\partial^2 \xi_2}{\partial i^2} + \frac{\partial^2 \xi_2}{\partial j^2} + \frac{\partial^2 \xi_2}{\partial k^2} \right) - \frac{1}{\rho} \frac{\partial p}{\partial j} \quad (2)$$
$$\frac{\partial \xi_3}{\partial s} + \xi_1 \frac{\partial \xi_3}{\partial i} + \xi_2 \frac{\partial \xi_3}{\partial j} + \xi_3 \frac{\partial \xi_3}{\partial k} = \frac{\mu}{\rho} \left( \frac{\partial^2 \xi_3}{\partial i^2} + \frac{\partial^2 \xi_3}{\partial j^2} + \frac{\partial^2 \xi_3}{\partial k^2} \right) - \frac{1}{\rho} \frac{\partial p}{\partial k}$$

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where

$$\frac{\partial \xi_1}{\partial i} + \frac{\partial \xi_2}{\partial j} + \frac{\partial \xi_3}{\partial k} = 0$$

Navier-Stokes equations should also follow a variational principle, though much effort has been paid, its existence is still a big problem. Scientists only find some variational formulations for simple fluids [4-6]. A variational principle can give profound, original, and challenging insights of a fluid problem, especially the travelling waves.

The well-known KdV equations [7-9] is the approximate case of the Navier-Stokes equations, there are various variational principles for KdV equations [10-12], and the modern soliton theory is originally developed from the KdV equation. This paper aims at searching for solitary waves directly from the Navier-Stokes equations by establishment of a suitable variational principle.

### Variational principle

The variational formulation for eq. (1) is extremely difficult to be obtained. This paper is to search for solitary waves from Navier-Stokes equations, so we focus ourselves on a constrained variational formulation by the following transformations [13-16]:

$$\begin{aligned}\xi_1(i, j, k, s) &= \Xi(\varepsilon) \\ \xi_2(i, j, k, s) &= A(\varepsilon) \\ \xi_3(i, j, k, s) &= B(\varepsilon) \\ p(i, j, k, s) &= P(\varepsilon)\end{aligned}\quad (3)$$

$$\varepsilon = \lambda_1 i + \lambda_2 j + \lambda_3 k - cs + \varepsilon_0 \quad (4)$$

$$\begin{aligned}\alpha \xi_2 &= \xi_1 \\ \beta \xi_3 &= \xi_1 \\ \vartheta p &= \xi_1\end{aligned}\quad (5)$$

where  $\alpha$ ,  $\beta$ , and  $\vartheta$  are all non-zero functions.

Based on the previous transformation, we can convert the Navier-Stokes equations into the following ordinary differential equation:

$$-c\Xi' + \lambda_1 \Xi \cdot \Xi' + \frac{\lambda_2}{\alpha} \Xi \cdot \Xi' + \frac{\lambda_3}{\beta} \Xi \cdot \Xi' = \frac{\mu \lambda_1^2}{\rho} \Xi'' + \frac{\mu \lambda_2^2}{\rho} \Xi'' + \frac{\mu \lambda_3^2}{\rho} \Xi'' - \frac{\lambda_1}{\rho \vartheta} \Xi' \quad (6)$$

Through the previous equation, we have:

$$-\frac{\mu}{\rho} (\lambda_1^2 + \lambda_2^2 + \lambda_3^2) \Xi'' + \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right) \Xi \cdot \Xi' + \left( \frac{\lambda_1}{\rho \vartheta} - c \right) \Xi' = 0 \quad (7)$$

Integrating (7), we have:

$$-\frac{\mu}{\rho} (\lambda_1^2 + \lambda_2^2 + \lambda_3^2) \Xi' + \frac{1}{2} \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right) \Xi^2 + \left( \frac{\lambda_1}{\rho \vartheta} - c \right) \Xi = H \quad (8)$$

According to (7) and (8), we have:

$$\Xi' = \frac{\rho}{\mu(\lambda_1^2 + \lambda_2^2 + \lambda_3^2)} \left[ \frac{1}{2} \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right) \Xi^2 + \left( \frac{\lambda_1}{\rho\vartheta} - c \right) \Xi - H \right] \quad (9)$$

Hence, eq. (7) can be represented:

$$\begin{aligned} & \frac{2\mu^2}{\rho} (\lambda_1^2 + \lambda_2^2 + \lambda_3^2)^2 \Xi'' - \rho \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right)^2 \Xi^3 - 3\rho \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right) \left( \frac{\lambda_1}{\rho\vartheta} - c \right) \Xi^2 \\ & \quad + \frac{-2\rho \left( \frac{\lambda_1}{\rho\vartheta} - c \right)^2 \Xi + 2\rho H \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right) \Xi + 2\rho H \left( \frac{\lambda_1}{\rho\vartheta} - c \right)}{2\mu(\lambda_1^2 + \lambda_2^2 + \lambda_3^2)} = 0 \end{aligned} \quad (10)$$

So the variational formulation of (10) can be established by the semi-inverse method [17], which is:

$$\begin{aligned} J(\Xi) = \int & \left[ -\frac{\mu}{2\rho} (\lambda_1^2 + \lambda_2^2 + \lambda_3^2) (\Xi')^2 - \frac{\rho \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right)^2}{12\mu(\lambda_1^2 + \lambda_2^2 + \lambda_3^2)} \Xi^4 - \right. \\ & \left. - \frac{\rho \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right) \left( \frac{\lambda_1}{\rho\vartheta} - c \right)}{2\mu(\lambda_1^2 + \lambda_2^2 + \lambda_3^2)} \Xi^3 \right] d\varepsilon + \\ & + \int \left[ -\rho \frac{\left( \frac{\lambda_1}{\rho\vartheta} - c \right)^2}{2\mu(\lambda_1^2 + \lambda_2^2 + \lambda_3^2)} \Xi^2 + \frac{\rho \left( \frac{\lambda_1}{\rho\vartheta} - c \right) H}{\mu(\lambda_1^2 + \lambda_2^2 + \lambda_3^2)} \Xi \right] d\varepsilon \end{aligned} \quad (11)$$

The semi-inverse method [17] has been widely used to establish a suitable variational formulation from a governing equation, for examples, the variational principle for singular waves [18], water waves [19], nano/micromechanical systems [20], two-point boundary value problems [20], KdV-Burgers-Kuramoto equation [21], 3D unsteady fluids [22], thin films [23], solitary waves [24], long water waves [25], Schrodinger equation [26].

From eq. (11), He-Weierstrass function [27] can be obtained:

$$E(\varepsilon, \Xi, \Xi', \omega) = \frac{1}{2} \omega^2 - \frac{1}{2} (\Xi')^2 - (\omega - \Xi') \Xi' \quad (12)$$

where  $\omega = \partial \Xi / \partial \varepsilon$ .

From eq. (12), It is evident that:

$$E(\varepsilon, \Xi, \Xi', \omega) = 0, \quad \frac{\partial^2 E}{\partial \omega^2} > 0 \quad (13)$$

Equation (13) shows that eq. (12) is a minimal variational principle.

### Solitary wave solutions

The objective of this section is to identify solitary wave solutions for Navier-Stokes equations by the obtained variational principle. The idea goes back to [28] variational approach to solitons, and it has been showing its validity for various wave equations, for examples, Boussinesq equation [29] and various the wave equations [30-36].

According to He's variational theory [28], we assume that the solitary solution of eq. (11) take the following form:

$$\Xi(\varepsilon) = \kappa \operatorname{sech}^2(v\varepsilon) \quad (14)$$

where  $\kappa \neq 0$ ,  $v \neq 0$ ,  $\kappa$  and  $v$  are unknown constants to be determined later.

Upon simultaneous solution of eq. (11) and eq. (14), the resulting expression is shown [28]:

$$\begin{aligned} J(\kappa, v) = & \int_0^\infty \left\{ -\frac{2\mu}{\rho} (\lambda_1^2 + \lambda_2^2 + \lambda_3^2) \kappa^2 v^2 [\operatorname{sech}^2(v\varepsilon) \operatorname{tg}(v\varepsilon)]^2 - \right. \\ & \left. - \frac{\rho \kappa^4}{12\mu(\lambda_1^2 + \lambda_2^2 + \lambda_3^2)} \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right)^2 [\operatorname{sech}^2(v\varepsilon)]^4 \right\} d\varepsilon - \\ & - \int_0^\infty \left\{ \frac{\rho \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right) \left( \frac{\lambda_1}{\rho g} - c \right) \kappa^3}{2\mu(\lambda_1^2 + \lambda_2^2 + \lambda_3^2)} [\operatorname{sech}^2(v\varepsilon)]^3 + \right. \\ & \left. + \frac{\left[ \left( \frac{\lambda_1}{\rho g} - c \right)^2 - \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right) H \right] \rho \kappa^2}{2\mu(\lambda_1^2 + \lambda_2^2 + \lambda_3^2)} [\operatorname{sech}^2(v\varepsilon)]^2 \right\} d\varepsilon + \\ & + \int_0^\infty \left\{ \frac{\rho H \kappa}{\mu(\lambda_1^2 + \lambda_2^2 + \lambda_3^2)} \left( \frac{\lambda_1}{\rho g} - c \right) \operatorname{sech}^2(v\varepsilon) \right\} d\varepsilon \end{aligned} \quad (15)$$

The following are the results obtained:

$$\begin{aligned} J(\kappa, v) = & - \frac{\rho \kappa 28 \frac{\mu^2}{\rho^2} (\lambda_1^2 + \lambda_2^2 + \lambda_3^2)^2 \kappa v^2}{105 v \mu (\lambda_1^2 + \lambda_2^2 + \lambda_3^2)} - \\ & - \frac{\rho \kappa \left[ 4 \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right)^2 \kappa^3 + 28 \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right) \left( \frac{\lambda_1}{\rho g} - c \right) \kappa^2 \right]}{105 v \mu (\lambda_1^2 + \lambda_2^2 + \lambda_3^2)} + \\ & + \frac{\rho \kappa \left[ -35 \left( \frac{\lambda_1}{\rho g} - c \right)^2 \kappa + 35 H \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right) \kappa + 105 \left( \frac{\lambda_1}{\rho g} - c \right) H \right]}{105 v \mu (\lambda_1^2 + \lambda_2^2 + \lambda_3^2)} \end{aligned} \quad (16)$$

Considering He's variational method [28], it gives:

$$\frac{\partial J}{\partial \kappa} = 0 \tag{17}$$

$$\frac{\partial J}{\partial v} = 0 \tag{18}$$

When

$$\left( \frac{\lambda_1}{\rho \vartheta} - c \right)$$

tends to zero, the eq. (16) can bring the following results:

$$\frac{-56 \frac{\mu^2}{\rho} (\lambda_1^2 + \lambda_2^2 + \lambda_3^2)^2 \kappa v^2 - 16 \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right)^2 \rho \kappa^3 + 70H \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right) \rho \kappa}{105v\mu(\lambda_1^2 + \lambda_2^2 + \lambda_3^2)} = 0 \tag{19}$$

$$\frac{-28 \frac{\mu^2}{\rho} (\lambda_1^2 + \lambda_2^2 + \lambda_3^2)^2 \kappa^2 v^2 + 4 \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right)^2 \rho \kappa^4 - 35H \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right) \rho \kappa^2}{105v^2\mu(\lambda_1^2 + \lambda_2^2 + \lambda_3^2)} = 0 \tag{20}$$

Solving eqs. (19) and (20) we can determine  $\kappa$  and  $v$ :

$$\kappa = \frac{\sqrt{35H}}{\sqrt{6 \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right)}} \tag{21}$$

$$v = \rho \sqrt{\frac{5 \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right)}{12H\mu^2(\lambda_1^2 + \lambda_2^2 + \lambda_3^2)^2}} \tag{22}$$

So the eq. (14) can be replaced by:

$$\Xi(\varepsilon) = \frac{\sqrt{35H}}{\sqrt{6 \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right)}} \operatorname{sech}^2 \left[ \rho \sqrt{\frac{5 \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right)}{12H\mu^2(\lambda_1^2 + \lambda_2^2 + \lambda_3^2)^2}} (\lambda_1 i + \lambda_2 j + \lambda_3 k - cs + \varepsilon_0) \right] \tag{23}$$

With eq. (11), the solitary wave solution of eq. (1) can be approximated:

$$\Xi(i, j, k, s) = \frac{\sqrt{35H}}{\sqrt{6 \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right)}} \operatorname{sech}^2 \cdot \left[ \rho \sqrt{\frac{5 \left( \lambda_1 + \frac{\lambda_2}{\alpha} + \frac{\lambda_3}{\beta} \right)}{12H\mu^2(\lambda_1^2 + \lambda_2^2 + \lambda_3^2)^2}} (\lambda_1 i + \lambda_2 j + \lambda_3 k - cs + \varepsilon_0) \right] \tag{24}$$

**Two examples**

This section gives two examples to verify the accuracy and effectiveness of He’s variational method [28].

*Example 1.* Consider variables in eq. (24), let  $\lambda_1 = 2, \lambda_2 = 1, \lambda_3 = 1, \alpha = 2, \beta = 2, \rho = 1, \vartheta = 1, c = 3, \mu = -1,$  and  $H = 1.$

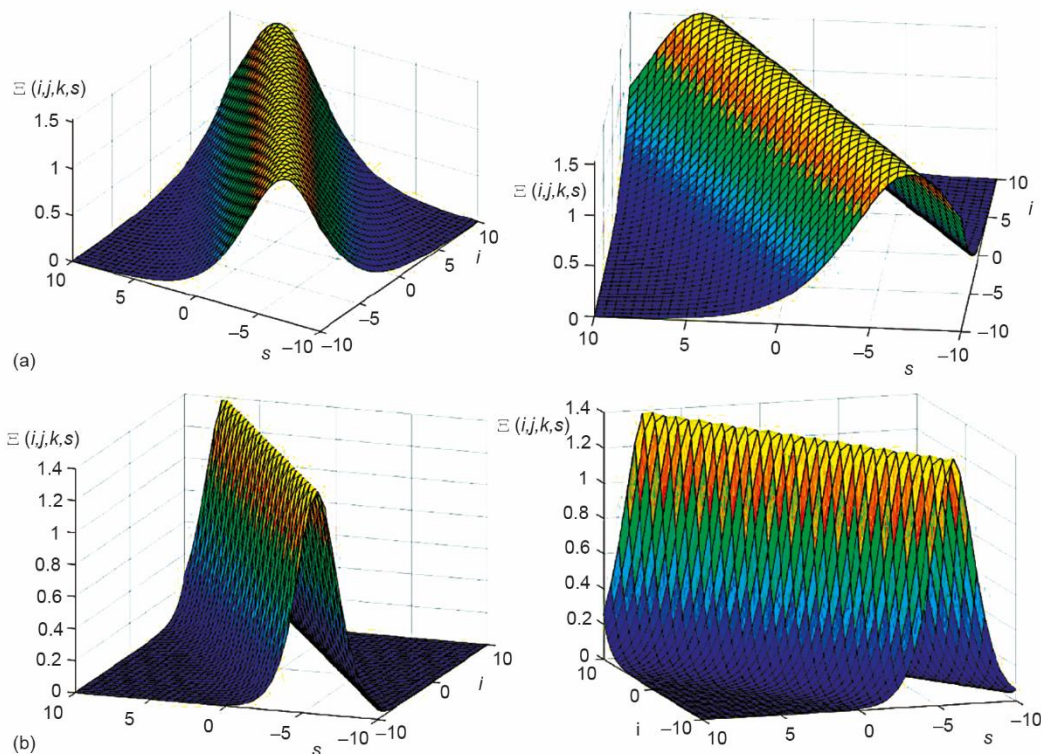
When  $j = 0, k = 0,$  and  $\varepsilon_0 = 1,$  we obtain the solitary wave solution, fig. 1(a) in a single direction:

$$\Xi(i, j, k, s) = \sqrt{\frac{35}{18}} \operatorname{sech}^2 \left[ \sqrt{\frac{5}{144}} (2i - 3s + 1) \right] \tag{25}$$

*Example 2.* In this case, we use  $\lambda_1 = 1, \lambda_2 = 2, \lambda_3 = 3, \alpha = 2, \beta = 1, \rho = -2, \vartheta = 1, c = -3, \rho = 2,$  and  $H = 2.$

When  $i = 0, k = 0,$  and  $\varepsilon_0 = -2,$  the solitary wave solution, fig. 1(b), in another single direction can be obtained:

$$\Xi(i, j, k, s) = \sqrt{\frac{7}{3}} \operatorname{sech}^2 \left[ \frac{5}{14} \sqrt{\frac{1}{24}} (2j + 3s - 2) \right] \tag{26}$$



**Figure 1. Solitary wave solutions for Navier-Stokes equations; (a)  $\lambda_1 = 2, \lambda_2 = 1, \lambda_3 = 1, \alpha = 2, \beta = 2, \rho = 1, \vartheta = 1, c = 3, \mu = -1,$  and  $H = 1,$  (b)  $\lambda_1 = 1, \lambda_2 = 2, \lambda_3 = 3, \alpha = 2, \beta = 1, \rho = -2, \vartheta = 1, c = -3, \mu = 2,$  and  $H = 2$  (for color image see journal web site)**

## Conclusions

Navier-Stokes millennium-prize problem is still an open problem [37], and there might have not exact solution to the Navier-Stokes equations, though the model has been widely used to explain various unsolved problems, *e.g.*, the mountain-river-desert relation [38, 39] and the dynamical properties of a rotating rigid body containing a viscous incompressible fluid [40]. This paper gives an alternative approach to the open problem, and exact solutions exist for solitary waves.

In this work, the solitary wave solution of Navier-Stokes equations has been obtained by He's variational method. This paper offers a totally new window for searching for solitary wave solutions directly from Navier-Stokes equations instead of its various approximate forms like KdV equation or Burgers equation, making the soliton theory much accurate to model the solitary waves.

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## References

- [1] Hill, E. L., Hamilton Principle and Conservation Theorem of Mathematical Physics, *Reviews of Modern Physics*, 23 (1951), 3, pp. 253-260
- [2] Ma, H. J., Simplified Hamiltonian-Based Frequency-Amplitude Formulation for Non-linear Vibration Systems, *Facta Universitatis Series: Mechanical Engineering*, 20 (2022), 2, pp. 445-455
- [3] He, J.-H., et al., Forced Non-Linear Oscillator in a Fractal Space, *Facta Universitatis Series: Mechanical Engineering*, 20 (2022), 1, pp. 1-20
- [4] Finlayson, B. A., Existence of Variational Principles for Navier-Stokes Equation, *Physics of Fluids*, 15 (1972), 6, pp. 963-967
- [5] Yasue, K., A Variational Principle for the Navier-Stokes Equation, *Journal of Functional Analysis*, 51 (1983), 2, pp. 133-141
- [6] Kerswell, R. R., Variational Principle for the Navier-Stokes Equations, *Physical Review E*, 59 (1999), 5, pp. 5482-5494
- [7] Kenig, C. E., et al., A Bilinear Estimate with Applications to the KdV Equation, *Journal of the American Mathematical Society*, 9 (1996), 2, pp. 573-603
- [8] He, J.-H., et al., Exp-Function Method for Non-linear Wave Equations, *Chaos Solitons & Fractals*, 30 (2006), 3, pp. 700-708
- [9] Cao, X. Q., et al., Variational Principles for Two Kinds of Non-Linear Geophysical Kdv Equation with Fractal Derivatives, *Thermal Science*, 26 (2022), 3B, pp. 2505-2515
- [10] Shen, Y., He, J. H., Variational Principle for a Generalized KdV Equation in a Fractal Space, *Fractals*, 28 (2020), 4, 2050069
- [11] He, J. H., Variational Principle for the Generalized KdV-Burgers Equation with Fractal derivatives for Shallow Water Waves, *J. Appl. Comput. Mech.*, 6 (2020), 4, pp. 735-740
- [12] Sun, J. S., Traveling Wave Solution of Fractal KDV-Burgers-Kuramoto Equation Within Local Fractional Differential Operator, *Fractals*, 29 (2021), 7, 2150231
- [13] Weekes, S. L., The Travelling Wave Scheme for The Navier-Stokes Equations, *SIAM. J. Numer. Anal.*, 35 (1998), 3, pp. 1249-1270
- [14] Dubovskii, C. P., et al., Travelling Wave-Like Solutions of the Navier-Stokes and the Related Equations, *J. Math. Anal. Appl.*, 204 (1996), 0477, pp. 930-939
- [15] Cazacu, C. A., et al., Transformation of The Travelling Wave Shape in Propagation on A Straight and Inclined Bed, *Stud. U. Babeş-Bol. Mat.*, 57 (2012), 2, pp. 167-173

- [16] Bakhoum, E. G., Cristian, T., Mathematical Transform of Traveling-Wave Equations and Phase Aspects of Quantum Interaction, *Math. Probl. Eng.*, 2010 (2010), 695208
- [17] He, J.-H., Variational Principles for Some Non-Linear Partial Differential Equations with Variable Coefficients, *Chaos Solitons & Fractals*, 19 (2004), 4, pp. 847-851
- [18] He, C. H., Liu, C., Variational Principle for Singular Waves, *Chaos, Solitons & Fractals*, 172 (2023), 113566
- [19] Wang, K. L., He, C. H., A Remark on Wang's Fractal Variational Principle, *Fractals*, 27 (2019), 1950134
- [20] He, C. H., A Variational Principle for a Fractal Nano/Microelectromechanical (N/MEMS) System, *Int. J. Numer. Methods H.*, 33 (2023), 1, pp. 351-359
- [21] Wang, S. Q., A Variational Approach to Non-Linear Two-Point Boundary Value Problems, *Computers & Mathematics with Applications*, 58 (2009), 11, pp. 2452-2455
- [22] He, J.-H., Lagrange Crisis and Generalized Variational Principle For 3D Unsteady Flow, *Int. J. Numer. Method. H.*, 30 (2019), 3, pp. 1189-1196
- [23] He, J. H., Sun C., A Variational Principle for a Thin Film Equation. *J. Math. Chem.*, 57 (2019), 9, pp. 2075-2081
- [24] Liu, M. Z., et al., Internal Solitary Waves in The Ocean by Semi-Inverse Variational Principle, *Thermal Science*, 26 (2022), 3B, pp. 2517-2525
- [25] Sun, J. S., Variational Principle for Fractal High-Order Long Water-Wave Equation, *Thermal Science*, 27 (2023), 3A, pp. 1899-1905
- [26] Sun, J. S., Fractal Modification of Schrodinger Equation and Its Fractal Variational Principle, *Therm. Sci.*, 27 (2023), accepted
- [27] He, J.-H., et al., On a Strong Minimum Condition of a Fractal Variational Principle, *Appl Math Lett*, 119 (2021), 107199
- [28] He, J. H., Asymptotic Methods for Solitary Solutions and Compactons, *Abstr. Appl. Anal.*, 2012 (2012), pp. 97-102
- [29] Wang, K. J., Wang, G. D., Solitary and Periodic Wave Solutions of The Generalized Fourth-Order Boussinesq Equation Via He's Variational Methods, *Math. Method. Appl. Sci.*, 44 (2021), 7, pp. 5617-5625
- [30] Wang, K. J., et al., Solitary Waves of The Fractal Regularized Long-Wave Equation Traveling Along an Unsmooth Boundary, *Fractals*, 30 (2022), 1, pp. 1-6
- [31] He, J.-H., et al., Solitary Waves Travelling Along an Unsmooth Boundary, *Results in Physics*, 24 (2021), 104104
- [32] He, J.-H., et al., Solitary Waves of The Variant Boussinesq-Burgers Equation in a Fractal-Dimensional Space, *Fractals*, 30 (2022), 3, 2250056
- [33] Sun, J. S., Approximate Analytic Solutions of Multi-Dimensional Fractional Heat-Like Models with Variable Coefficients, *Thermal Science*, 23 (2019), 6B, pp. 3725-3729
- [34] Wang, K. L., et al., A Novel Perspective to the Local Fractional Bidirectional Wave Model on Cantor Sets, *Fractals*, 30 (2022), 6, pp. 1-7
- [35] Sun, J. S., Variational Principle and Solitary Wave of the Fractal Fourth-Order Non-linear Ablowitz-Kaup-Newell-Segur Water Wave Model, *Fractals*, 31 (2023), 5, 2350036
- [36] Sun, J.S., Variational Principle for Fractal High-Order Long Water-Wave Equation, *Thermal Science*, 27 (2023), 3, pp. 1899-1905
- [37] Nelkin, M., In What Sense is Turbulence an Unsolved Problem, *Science*, 255 (1992), Jan., pp. 566-570
- [38] Mei, Y., et al., The Yellow River-Bed Evolution: A Statistical Proof of the Mountain-River-Desert Conjecture, *Thermal Science*, 27 (2023), 3A, pp. 2075-2079
- [39] Mei, Y., et al., On the Mountain-River-Desert Relation, *Thermal Science*, 25 (2021), 6, pp. 4817-4822
- [40] He, J.-H., et al., Dynamical Analysis of a Rotating Rigid Body Containing a Viscous Incompressible Fluid, *International Journal of Numerical Methods for Heat & Fluid Flow*, 33 (2023), 8, pp. 2800-2814