



Parameter uniform fnite diference formulation with oscillation free for solving singularly perturbed delay parabolic diferential equation via exponential spline

Zerihun Ibrahim Hassen<sup>1\*†</sup> and Gemechis File Duressa<sup>2†</sup>

# **Abstract**

**Objective** In this work, singularly perturbed time dependent delay parabolic convection-difusion problem with Dirichlet boundary conditions is considered. The solution of this problem exhibits boundary layer at the right of special domain. In this layer the solution experiences steep gradients or oscillation so that traditional numerical methods may fail to provide smooth solutions. We developed oscillation free parameter uniform exponentially spline numerical method to solve the considered problem.

**Results** In the temporal direction, the implicit Euler method is applied, and in the spatial direction, an exponential spline method with uniform mesh is applied. To handle the effect of perturbation parameter, an exponential fitting factor is introduced. For the developed numerical scheme, stability and uniform error estimates are examined. It is shown that the scheme is uniformly convergent of linear order in the maximum norm. Numerical examples are provided to illustrate the theoretical fndings.

**Keywords** Exponential spline, Oscillation-free, Singularly perturbed delay problem, Fitting factor, Convectiondifusion, Uniform convergence

**Mathematics Subject Classifcation** 65M06, 65M12, 65M15

# **Introduction**

Delay diferential equations (DDEs) are a class of differential equation where the unknown function or its derivative at a certain time depends on the solution and possibly its derivatives at earlier times. The delay in these equations represents the transport delay, incubation

† The authors Zerihun Ibrahim Hassen and Gemechis File Duressa have contributed equally to this work.

\*Correspondence:

Zerihun Ibrahim Hassen

zerihunibrahim@gmail.com

<sup>1</sup> Department of Mathematics, Arba Minch University, Arba Minch, Ethiopia

<sup>2</sup> Department of Mathematics, Jimma University, Jimma, Ethiopia

period, gestation time, etc. If a small positive parameter  $\varepsilon$  multiply the highest derivative term of DDEs and involves at least a delay term, the DDEs are said to be singularly perturbed delay diferential equations (SPD-DEs). When the delay parameter magnitude is larger than the perturbation parameter, the equations are said to SPDDEs with large delay, otherwise they are said to be SPDDEs with small delay. These problems can be found various applications in science and engineering such as control systems [[1\]](#page-13-0), chemical reactions [\[2](#page-13-1)], epidemiology [[3\]](#page-13-2), optics and physiology [\[4](#page-13-3)], and neural networks [\[5](#page-13-4)].

Singularly perturbed delay parabolic convectiondifusion problems (SPDPCDPs) are a type of SPDDE. SPDPCDPs are signifcant in modeling various physical phenomena, particularly in systems where the current



© The Author(s) 2025 **Open Access** This article is licensed under a Creative Commons Attribution-NonCommercial-NoDerivatives 4.0 International License, which permits any non-commercial use, sharing, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons licence, and indicate if you modifed the licensed material. You do not have permission under this licence to share adapted material derived from this article or parts of it. The images or other third party material in this article are included in the article's Creative Commons licence, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons licence and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this licence, visit<http://creativecommons.org/licenses/by-nc-nd/4.0/>.

state infuences future behavior, such as in fuid dynamics and heat transfer. These problems are crucial in modeling systems where there is a signifcant disparity between the rates of convection, difusion, and delay efects. For instance, in chemical engineering [\[2](#page-13-1)], they can describe the behavior of reactive transport in porous media where the reaction rates and transport processes difer vastly, and delays in the system (due to transport lags or reaction time) afect the overall dynamics. Analyzing these problems helps in understanding how the interplay between fast and slow processes, coupled with delay efects, infuences the stability and evolution of the system.

The solution for SPDDEs has a boundary layer because of the perturbation parameter  $\varepsilon$ . Thus, the solution shows large variation, oscillation, in small region of the domain. It is long familiar that most classical numerical methods are unable to give accurate result on uniform mesh for such problems especially as  $\varepsilon$  goes to zero unless a very fne mesh is considered, which is computationally expensive. Hence there is the need for methods which are stable and uniformly convergent irrespective of the values of  $\varepsilon$  and mesh size [[6\]](#page-13-5). Finite difference numerical methods that exhibit uniform convergence and stable are mainly developed using ftted meshes and ftted operators. While ftted operator methods maintain a uniform mesh, fttedmesh methods concentrate on selecting a fne mesh in the layer region(s). There are also other nonclassical finite diference numerical methods to solve singularly perturbed delay diferential equations such as adaptive mesh refnement and domain decomposition methods.

SPDPCDPs are studied by various authors. Kaushik and Sharma [\[7](#page-13-6)] approximated SPDPCDPs using a weighted diference time discretization and central diference space discretization on a piecewise Shishkin mesh. They have shown that the method is stable and uniformly convergent with respect to  $\varepsilon$ . SPDPCDPs are estimated by Das and Natesan [\[8](#page-13-7)] using for time derivative an implicit-Euler scheme and for spatial derivatives a hybrid scheme which made up of midpoint upwind scheme and the central difference scheme. To solve SPDPCDPs, Gowrisankar and Natesan [[9](#page-13-8)] used the upwind fnite diference scheme for spatial derivatives and the backward-Euler scheme for time derivatives. They proved the proposed method is parameter uniform convergent of frst order. Using the exponentially ftted operator fnite diference method for spatial discretization and the Crank-Nicolson method for temporal discretization, Woldaregay et al. [\[10\]](#page-13-9) solved SPDPCDPs through both approaches. They have shown that the proposed scheme converges uniformly with frst order of convergence. Negero and Duressa [[11](#page-13-10)] studied second order convergent scheme to approximate SPDPCDPs. After a year the authors [\[12\]](#page-13-11) constructed a second order accurate scheme for solving SPDPCDPs. Negero and Duressa [[13](#page-13-12)]

estimated the solution of SPDPCDPs using Crank-Nicolson's time discretization scheme and exponentially ftted cubic spline scheme for spatial discretization. Recently the following authors have developed parameter uniform convergent numerical scheme to solve SPDPCDPs. Hassen and Duressa [[14](#page-13-13)] approximated SPDPCDPs using Crank-Nicolson time discretization and upwind fnite diference for spatial derivative using Peano kernel theorem convergent analysis. Fitted computational method is developed by Tesfaye et al. [[15](#page-13-14)] to solve SPDPCDPs. After a year the authors [\[16\]](#page-13-15) solved SPDPCDPs by employing backward Euler scheme for time derivatives. They used a higher-order fnite diference method to approximate the second-order derivative and non-symmetric fnite diference schemes to approximate the frst-order derivative terms. Hassen and Duressa [\[17\]](#page-13-16) developed a parameter uniform convergent numerical scheme to solve SPDPCDPs by employing implicit Euler approach in the time direction and extended cubic B-spline collocation in the space direction. Kumar and Gowrisankar  $[18]$  have suggested an efficient numerical method for SPDPCDPs. The authors proved that the proposed numerical method converges uniformly with frst-order up to logarithm in the spatial variable and also frst-order in the temporal variable. Readers can refer different numerical scheme for solving SPDDEs in [[19](#page-13-18)[–32](#page-13-19)].

In this paper, our aim is to develop a parameter uniform numerical scheme for SPDPCDPs large delay version with Dirichlet boundary conditions. The proposed scheme comprises of implicit Euler in temporal direction and exponential spline scheme in spatial direction. We provided an exponentially ftting factor to manage the perturbation parameter's effects. The novelty of the presented scheme is that, unlike Shishkin and Bakhvalov mesh types, it does not depend on a specially designed mesh and needs no prior knowledge regarding the boundary layer's width and position. Results from the suggested scheme are more precise, consistent, and uniformly convergent.

*Notation* The symbols  $N_t$  and  $N_z$  are denoted for the number of mesh elements, mesh parameters, in time and space direction, respectively; the symbol *C* is denoted for a generic positive constant which is independent of perturbation parameter and mesh parameters. The norm  $\|.\|$  denotes supremum or maximum norm, i.e.,  $\|\pi(z, t)\| = \max_{(z, t) \in \Omega} |\pi(z, t)|$ .

## **Continuous problem**

Let  $\Omega_z = (0, 1), \Omega_t = (0, T]$  be are spatial and temporal domain respectively, and  $\Omega = \Omega_z \times \Omega_t$  for  $T > 0$ . We consider the following SPDPCDP of the form:

$$
\begin{cases}\nf_t(z,t) + \mathcal{L}_{\varepsilon}f(z,t) = -c(z,t)f(z,t-\kappa) + g(z,t), & (z,t) \in \Omega, \\
f(z,t) = B_b(z,t), & (z,t) \in \varphi_b = [0,1] \times [-\kappa,0], \\
f(0,t) = B_l(t), & t \in \varphi_l = \{(0,t) : 0 \le t \le T\}, \\
f(1,t) = B_r(t), & t \in \varphi_r = \{(1,t) : 0 \le t \le T\},\n\end{cases}
$$
\n(1)

where  $\mathcal{L}_{\varepsilon} f(z, t) = -\varepsilon f_{zz}(z, t) + a(z) f_z(z, t) + b(z, t) f(z, t)$ . The delay parameter in the given problem is denoted by  $\kappa > 0$ , and the perturbation parameter by  $\varepsilon \in (0,1]$ . The functions  $a(z)$ ,  $b(z, t)$ ,  $c(z, t)$ , and  $g(z, t)$  on  $\overline{\Omega} = [0, 1] \times [0, T]$ and  $B_b(z, t)$ ,  $B_l(t)$ , and  $B_r(t)$  on  $\varphi = \varphi_l \cup \varphi_b \cup \varphi_r$  are assumed sufficiently smooth and bounded which satisfy  $a(z) \ge \gamma > 0$ ,  $b(z, t) \ge \eta$  and  $c(z, t) \ge \beta$ . Assume T satisfy  $T = \zeta \kappa$ ,  $\zeta$  is positive integer. Under these circumstances the problem exhibits a boundary layer at the right side of the spatial domain.

The Hölder continuous of the data, together with the compatibility condition at the corner points delay term [[33\]](#page-13-20) as stated below, can ensure the existence and uniqueness of the solution to the problem ([1\)](#page-2-0).

$$
B_l(0) = B_b(0,0), \qquad B_r(0) = B_b(1,0), \tag{2}
$$

(3)  $\frac{dB_l(0)}{dt} - \varepsilon \frac{\partial^2 B_b(0,0)}{\partial z^2} + a(0) \frac{\partial B_b(0,0)}{\partial z} + b(0,0)B_b(0,0) = -c(0,0)B_b(0,-\kappa) + g(0,0),$  $\frac{dB_l(0)}{dt} - \varepsilon \frac{\partial^2 B_b(1,0)}{\partial z^2} + a(1) \frac{\partial B_b(1,0)}{\partial z} + b(1,0)B_b(1,0) = -c(1,0)B_b(1,-\kappa) + g(1,0).$ 

Let  $\mathcal{L} f(z,t) = f_t(z,t) + \mathcal{L}_{\varepsilon} f(z,t)$ , then the differential operator  $L$  satisfies the next Lemma.

<span id="page-2-2"></span>**Lemma 1** (Continuous Maximum Principle) *Let*   $y(z, t) \in C^2(\Omega) \cap C^0(\overline{\Omega})$ *, satisfies*  $\mathscr{L} y(z, t) > 0$  for all  $(z, t) \in \Omega$  and  $y(z, t) \ge 0$  for all  $(z, t) \in \varphi$ . Then  $y(z, t) \ge 0$ , *for all*  $(z, t) \in \overline{\Omega}$ .

*Proof* Let  $(z^*, t^*) \in \overline{\Omega}$ , such that  $y(z^*, t^*) = \min_z y(z, t)$  $(z,t)\in\overline{\Omega}$ and suppose that  $y(z^*, t^*)$  < 0. Obviously  $(z^*, t^*) \notin \varphi$  and  $(z^*, t^*) \in \Omega$ . From calculus property, we have  $y_z(z^*, t^*) = 0$ ,  $y_t(z^*, t^*) = 0$ , and  $y_{zz}(z^*, t^*) > 0$ . Hence from  $(1)$  $(1)$ , we have

$$
\mathcal{L} y(z^*, t^*) = y_t(z^*, t^*) - \varepsilon y_{zz}(z^*, t^*)
$$
  
+  $a(z^*) y_z(z^*, t^*) + b(z^*, t^*) y(z^*, t^*) < 0.$ 

<span id="page-2-0"></span>This contradicts the hypothesis  $\mathcal{L}v(z,t) > 0$ . Therefore  $y(z, t) > 0$  for all  $(z, t) \in \overline{\Omega}$ .

<span id="page-2-1"></span>**Lemma 2** *The solution*  $y(z, t)$  *of the problem* [\(1](#page-2-0)) *satisfies* 

$$
\left|f(z,t)-B_b(z,0)\right|\leq Ct,\quad (z,t)\in\overline{\Omega},
$$

*where a constant*  $C > 0$ , *does not depend on*  $\varepsilon$ *.* 

**Proof** Refer [8]. 
$$
\square
$$

**Lemma 3** *The solution*  $f(z, t)$  *of the problem* [\(1](#page-2-0)) *satisfies*  $|f(z,t)| \leq C$ ,  $(z,t) \in \overline{\Omega}$ , where a constant  $C > 0$ , does not *depend on* ε.

*Proof* From Lemma [2,](#page-2-1) we have

$$
|f(z,t)| \le |f(z,t) - B_b(z,0)| + |B_b(z,0)|
$$
  
\n
$$
\le Ct + |B_b(z,0)|
$$
  
\n
$$
\le C, \text{ since } t \in (0,T] \text{ and } B_b(z,0) \in C^2(\overline{\Omega}).
$$

This completes the proof.  $\Box$ 

The stability of the continuous differential operator  $L$ and an  $\varepsilon$ -uniform bound for the problem ([1](#page-2-0)) in the maximum norm are provided by the following Lemma. The Lemma result follows from the maximum principle.

Lemma 4 (Stability result for Continuous Problem) The *solution f*(*z*, *t*) *of* ([1](#page-2-0)) *satisfes*

$$
|f(z,t)| \leq \eta^{-1} ||\mathscr{L}f(z,t)|| + \max\{|B_l(t), B_r(t), B_b(z,t)|\}.
$$

*Proof* Let  $K = \max\{|B_l(t), B_r(t), B_b(z, t)|\}$ . For the barrier function  $\Lambda^{\pm}(z,t) = \eta^{-1} || \mathcal{L}f(z,t) || + K \pm f(z,t)$ , we have

$$
\Lambda^{\pm}(0,t) = \eta^{-1} || \mathcal{L}f(z,t) || + \max\{|B_l(t), B_r(t), B_b(0,t)|\} \pm f(0,t) > 0,\Lambda^{\pm}(1,t) = \eta^{-1} || \mathcal{L}f(z,t) || + \max\{|B_l(t), B_r(t), B_b(1,t)|\} \pm f(1,t) > 0,\mathcal{L} \Lambda^{\pm}(z,t) = b(z,t)[\eta^{-1} || \mathcal{L}f(z,t) || + K] \pm \mathcal{L}f(z,t)\ge || \mathcal{L}f(z,t) || + b(z,t)K \pm \mathcal{L}f(z,t)\ge || \mathcal{L}f(z,t) || \pm \mathcal{L}f(z,t)\ge 0.
$$

Hence by applying the Lemma  $1$ , we get the required  $r$ esult.

Furthermore, bounds on the solution and its derivatives are provided by the following theorem.

**Theorem 1** *The solution*  $f(z, t)$  *to the problem* [\(1](#page-2-0)) *and its derivatives satisfes*

$$
\left|\frac{\partial^{j+j}f(z,t)}{\partial z^j\partial t^i}\right|\leq C\Big(1+\varepsilon^{-j}\exp(-\gamma(1-z)/\varepsilon)\Big),\qquad (z,t)\in\overline{\Omega},
$$

*where i and j are non*-*negative integers such that*   $0 \le i + j \le 5$ .

**Proof** Refer on [34] 
$$
\square
$$

## **Numerical scheme**

## **Time discretization**

We engage a uniform mesh on the time domain [0, T] with time step size  $\Delta t$  as  $\Omega_t^{N_t} = \{ t_m = m\Delta t : m = 0(1)N_t, t_{N_t} = T, \ \Delta t = T/N_t \}$ and  $\Omega_t^P = \{t_s = -s\Delta t : s = 0(1)P, t_P = -\kappa\}$ , where  $N_t$  and *P* are th number of mesh elements in [0,  $T$ ] and  $[-\kappa, 0]$  respectively. In order to handle the term with delay, a special mesh is selected so that it coincides with a mesh point in  $\Omega_t^P$ . We use the implicit Euler scheme for time derivatives, so we obtain

$$
\frac{F^{m+1}(z) - F^m(z)}{\Delta t} - \varepsilon \frac{d^2 F^{m+1}(z)}{dz^2} + a(z) \frac{dF^{m+1}(z)}{dz} + b^{m+1}(z) F^{m+1}(z) = -c^{m+1}(z) F^{m+1-P}(z) + g^{m+1}(z)
$$
(4)

Consistently, we write ([4\)](#page-3-0) which gives semi-discrete scheme as

<span id="page-3-1"></span>
$$
\begin{cases}\n(I + \Delta t \mathcal{L}_\varepsilon^{* \Delta t}) F^{m+1}(z) = H^m(z), \\
F^{m+1}(0) = B_l(t_{m+1}), \quad F^{m+1}(1) = B_r(t_{m+1}), \\
F^{-s}(z, t) = B_b(z, -t_s), s = 0(1)P, z \in \overline{\Omega_z},\n\end{cases}
$$
\n(5)

where  $H^m(z) = F^m(z) + \Delta t \left( -c^{m+1}(z)F^{m+1-P}(z) + g^{m+1}(z) \right),$  $\mathscr{L}_{\varepsilon}^{* \Delta t} = -\varepsilon \frac{d^2}{dz^2} + a(z) \frac{d}{dz} + b^{m+1}(z)$ , and  $F^m(z)$  is the approximation of  $f(z, t)$  at  $t = t_m = m\Delta t$ , i.e.,  $F^m(z) \approx f(z, t_m)$ .

In [\(5\)](#page-3-1) let  $\mathscr{L}_{\varepsilon}^{\Delta t} = I + \Delta t \mathscr{L}_{\varepsilon}^{*\Delta t}$ . Then  $\mathscr{L}_{\varepsilon}^{\Delta t}$  satisfies the next maximum principle.

<span id="page-3-2"></span>**Lemma 5** (Semi-discrete Maximum Principle) *Let*  $\mu(z, t_{m+1}) \in C^2(\overline{\Omega_z})$ . *Assume*  $\mu(0, t_{m+1}) \geq 0$ ,  $\mu(1, t_{m+1}) \geq 0$ , and  $\mathcal{L}_{\varepsilon}^{\Delta t} \mu(z, t_{m+1}) \geq 0$  for all  $z \in \Omega_z$ , *then*  $\mu(z, t_{m+1}) \geq 0$  *for all*  $\overline{\Omega_z}$ *.* 

*Proof* Let  $(z^*, t_{m+1}) \in \{(z, t_{m+1}) : z \in \overline{\Omega_z}\}\)$  and  $\min_{m} \mu(z, t_{m+1}) = \mu(z^*, t_{m+1}) < 0.$  Clearly,  $z\in\overline{\Omega}$  $(z^*, t_{m+1}) \notin \{(0, t_{m+1}), (1, t_{m+1})\}.$  Also we have,  $\frac{d\mu(z^*,t_{m+1})}{dz} = 0$ ,  $\frac{d\mu(z^*,t_{m+1})}{dt} = 0$ , and  $\frac{d^2\mu(z^*,t_{m+1})}{dz^2} \ge 0$ . Then  $\mathcal{L}^{\Delta t}_{\varepsilon}\mu(z^*,t_{m+1}) = \frac{d\mu(z^*,t_{m+1})}{dt} + \Delta t \left(-\varepsilon\frac{d^2\mu(z^*,t_{m+1})}{dz^2}\right.$  $\overline{dz^2}$  $+a(z^*)\frac{d\mu(z^*,t_{m+1})}{dt}$ dz

This contradicts to the assumption made. Thus  $\mu(z^*, t_{m+1}) \geq 0$ , and hence  $\mu(z, t_{m+1}) \geq 0$  for all  $z \in \overline{\Omega_z}$ .  $\Box$ 

 $+ b^{m+1}(z^*) \mu(z^*, t_{m+1}) > 0.$ 

<span id="page-3-0"></span>The local truncation error  $e_{m+1}$  for the scheme [\(5](#page-3-1)) is given by  $e_{m+1} = f(z, t_{m+1}) - \bar{F}^{m+1}(z)$ , where  $\bar{F}^{m+1}(z)$  is the solution of

$$
\begin{cases}\n\left(I + \Delta t \mathcal{L}_\varepsilon^{* \Delta t}\right) \overline{F}^{m+1}(z) = f(z, t_m) + \Delta t \left(-c(z, t_{m+1})f(z, t_{m+1}) + g(z, t_{m+1})\right), \\
\overline{F}^{m+1}(0) = B_l(t_{m+1}), \quad \overline{F}^{m+1}(1) = B_r(t_{m+1}), \\
\overline{F}^{-s} = B_b(z, -t_s), \quad s = 0(1)P, z \in \overline{\Omega_z}.\n\end{cases} \tag{6}
$$

The local error in the time direction is estimated in the following lemma.

<span id="page-4-2"></span>**Lemma 6** (Local error) The local error  $e_{m+1}$  at  $t_{m+1}$  associated to the scheme  $(5)$  $(5)$  satisfies the bound  $||e_{m+1}|| \leq C(\Delta t)^2$ .

*Proof* The function  $\overline{F}^{m+1}(z)$  satisfies

$$
\begin{aligned} &\left(I + \Delta t \mathcal{L}_{\varepsilon}^{* \Delta t}\right) \overline{F}^{m+1}(z) \\ &- \Delta t \left(-c(z, t_{m+1})f(z, t_{m+1}) + g(z, t_{m+1})\right) = f(z, t_m). \end{aligned} \tag{7}
$$

From Taylor series expansion, we get

$$
f(z, t_m) = f(z, t_{m+1}) - \Delta t f_t(z, t_{m+1}) + \mathcal{O}((\Delta t)^2),
$$
  
\n
$$
= f(z, t_{m+1}) - \Delta t [-\mathcal{L}_{\varepsilon}^{*\Delta t} f(z, t_{m+1}) - c(z, t_{m+1})f(z, t_{m+1})] + \mathcal{O}((\Delta t)^2)
$$
  
\n
$$
= (I + \Delta t \mathcal{L}_{\varepsilon}^{*\Delta t}) f(z, t_{m+1}) + \mathcal{O}((\Delta t)^2).
$$
  
\n
$$
+ c(z, t_{m+1}) f(z, t_{m+1-p}) + g(z, t_{m+1}) + \mathcal{O}((\Delta t)^2).
$$
  
\n(8)

From [\(7](#page-4-0)) and [\(8](#page-4-1)) one can observe that  $e_{m+1}$  is the solution of

$$
\begin{cases}\n(I + \Delta t \mathcal{L}_{\varepsilon}^{* \Delta t}) e_{m+1} = \mathcal{O}((\Delta t)^2), \\
e_{m+1}(0) = 0 = e_{m+1}(1).\n\end{cases}
$$
\n(9)

Clearly the operator  $I + \Delta t \mathcal{L}_{\varepsilon}^{*\Delta t}$  satisfies semi-discrete maximum principle. Thus we obtain  $||e_{m+1}|| \leq C(\Delta t)^2$ .  $\Box$ 

 $E_{m+1} = f(z, t_{m+1}) - F^{m+1}(z)$  defines the global error in time direction by providing the error contribution at each time step.

<span id="page-4-3"></span>**Lemma** 7 (Global Error) *The global error*  $E_{m+1}$  at  $t_{m+1}$  associated to [\(5](#page-3-1)) satisfies  $\leq$   $||E_{m+1}|| \leq C \Delta t$ ,  $m = 0(1)N_t - 1$ .

*Proof* Using Lemma [6](#page-4-2), we get

$$
||E_{m+1}|| = \left\| \sum_{t=0}^{m+1} e_{t+1} \right\|
$$
  
\n
$$
\leq ||e_1|| + ||e_2|| + \cdots + ||e_{m+1}||
$$
  
\n
$$
\leq (m+1)C(\Delta t)^2
$$
  
\n
$$
= C((m+1)\Delta t)\Delta t, \quad (m+1)\Delta t \leq T
$$
  
\n
$$
\leq C\Delta t.
$$

As a result, the temporal discretization process is frst order uniform convergent.

The Lemmas  $5,6,$  $5,6,$  $5,6,$  and  $7$  show the stability and consistency of the scheme  $(5)$  $(5)$ . The derivative bound of the solution utilized to demonstrate the convergence of the method is provided by the following theorem.

<span id="page-4-6"></span>**Theorem 2** *The solution*  $F^{m+1}(z)$  *of* [\(5\)](#page-3-1) *satisfies the estimate*

<span id="page-4-0"></span>
$$
\left| \frac{d^{j}F^{m+1}(z)}{dz^{j}} \right| \leq C\Big(1 + \varepsilon^{-j}\exp(-\gamma(1-z)/\varepsilon)\Big),
$$
  
\n $j = 0(1)4, \quad z \in \overline{\Omega_{z}}$   
\n**Proof** Refer [16, 35]

## **Space discretization**

<span id="page-4-1"></span>We divide  $\overline{\Omega_z} = [0, 1]$  in to  $N_z$  equal number of subdomain with length of  $h = 1/N_z$  as  $0 = z_0, z_1, ..., z_{N_z} = 1$  and  $z_n = nh$ ,  $n = 0(1)N_z$ . Define  $\Omega_z^{N_z} = \{z_n = nh, n = 0(1)N_z\}$ and  $\Omega^N = \Omega_z^{N_z} \times \Omega_t^{N_t}$ . At  $z_n = nh$  for  $(m + 1)^{th}$  time level, let  $F_n^{m+1}$  be an approximation to  $F^{m+1}(z_n)$  found by exponential spline function  $Q_n(z)$  passing through the points  $(z_n, F_n^{m+1})$  and  $(z_{n+1}, F_{n+1}^{m+1})$ . Omit the superscript  $m + 1$ for convenience, i.e., $F_n^{m+1} = F_n$ . For each  $n^{th}$  segment, the exponential spline function has the following form [[36](#page-13-23)]:

<span id="page-4-5"></span><span id="page-4-4"></span>
$$
Q_n(z) = \phi_n \exp(\lambda(z - z_n)) + \chi_n \exp(-\lambda(z - z_n))
$$
  
+  $\psi_n(z - z_n) + \vartheta_n$ ,  $n = 1(1)N_z - 1$ , (10)

where  $\phi_n$ ,  $\chi_n$ ,  $\psi_n$ , and  $\vartheta_n$  are constants to be determined and  $\lambda$  is a free parameter used to advance the accuracy of the scheme. Here  $Q_n(z) \in C^2[\overline{\Omega_z}]$  interpolate  $F_n$  at  $z_n$ ,  $n = 0(1)N_z$  depends on  $\lambda$  and reduce to cubic spline  $Q_n(z)$  in  $\overline{\Omega_z}$  as  $\lambda \to 0$ . To obtain the coefficients intro-duced in ([10](#page-4-4)),  $Q_n(z)$  should satisfy the condition of first derivative continuity at the common nodes. Defne

$$
Q_n(z_n) = F_n, \quad Q_n(z_{n+1}) = F_{n+1},
$$
  
\n
$$
Q_n^{''}(z_n) = M_n, \quad Q_n^{''}(z_{n+1}) = M_{n+1}
$$
\n(11)

From [\(9](#page-4-5)) and [\(10](#page-4-4)) after some manipulation, we get

$$
\phi_n = h^2 \frac{M_{n+1} - \exp(-\xi)M_n}{2\xi^2 \sinh(\xi)}, \quad \chi_n = h^2 \frac{\exp(\xi)M_n - M_{n+1}}{2\xi^2 \sinh(\xi)}
$$

$$
\psi_n = \frac{F_{n+1} - F_n}{h} - h \frac{M_{n+1} - M_n}{\xi^2}, \quad \vartheta_n = F_n - h^2 \frac{M_n}{\xi^2},
$$
(12)

where  $\xi = \lambda h$  and  $n = 1(1)N_z - 1$ . From the first derivative continuity  $Q'_{n-1}(z_n) = Q'_{n}(z_n)$  for  $n = 1(1)N_z - 1$ , we obtain the following relations:

where  $L_1 = \frac{\sinh(\xi) - \xi}{\xi^2 \sinh(\xi)}, \quad L_2 = \frac{\xi \cosh(\xi) - \sinh(\xi)}{\xi^2 \sinh(\xi)}, \quad \text{and}$  $M_{\nu} = F^{''}(z_{\nu}), \nu = n, n \pm 1$ . Equation [\(4](#page-3-0)) can be written as

$$
-\varepsilon F^{''}(z) = a(z)F^{'}(z) + p(z)F(z) - G(z),
$$
 (14)

where  $F(z) = F^{m+1}(z)$ ,  $p(z) = p^{m+1}(z) = \frac{1}{\Delta t} + b^{m+1}(z)$ ,  $G(z) = G^m(z) = \frac{1}{\Delta t} F^m(z) - c^{m+1}(z) F^{m+1-P}(z) + g^{m+1}(z)$ . Equation [\(14](#page-5-0)) discretized by the exponential spline by introducing a fitting factor  $\sigma(\rho)$ , we get

$$
\varepsilon\sigma(\rho)M_{\nu}=a_{\nu}F_{\nu}^{'}+p_{\nu}F_{\nu}-G_{\nu},\quad\text{for }\nu=n,n\pm1,\tag{15}
$$

where  $\rho = h/\varepsilon$ . Placing exact solution in [\(15\)](#page-5-1) and substitute the resulting in to  $(13)$  $(13)$ , we obtain

<span id="page-5-6"></span>
$$
TE(h) = \frac{h^4}{3} (L_2 - 2L_1) a_n F^{(3)}(\varrho_n)
$$
  
+  $\varepsilon \sigma(\rho) \frac{h^4}{12} (-12L_1 + 1) a_n F^{(4)}(\varrho_n) + \mathcal{O}(h^6).$  (17)

<span id="page-5-2"></span><span id="page-5-0"></span>Clearly  $TE(h) = \mathcal{O}(h^4)$  for  $L_1 + L_2 = 1/2$ .<br>If  $L_2 = 5/12, L_1 = 1/12$ , we have If  $L_2 = 5/12, L_1 = 1/12$ , we have  $TE(h) = \varepsilon \sigma(\rho) \frac{h^6}{240} F^{(6)}(\rho_n), \rho_n \in [z_{n-1}, z_{n+1}].$  Using nonsymmetric fnite diference approximation for frst deriv-ative [\[38](#page-13-25)]  $F(z_n)$ , we have

<span id="page-5-3"></span>
$$
F^{'}(z_{n}) = \frac{F(z_{n+1}) - F(z_{n-1})}{2h} + \mathcal{O}(h^{2}),
$$
  
\n
$$
F^{'}(z_{n-1}) = \frac{-F(z_{n+1}) + 4F(z_{n}) - 3F(z_{n-1})}{2h} + \mathcal{O}(h^{2}),
$$
  
\n
$$
F^{'}(z_{n+1}) = \frac{3F(z_{n+1}) - 4F(z_{n}) + F(z_{n-1})}{2h} + \mathcal{O}(h^{2}).
$$
\n(18)

<span id="page-5-1"></span>Finally by substituting the approximation of ([18\)](#page-5-3) in to the approximation of  $(16)$  $(16)$ , we get the following full discretized scheme:

$$
\begin{cases}\n\mathcal{L}_{\varepsilon}^{\Delta t,h} F_{n}^{m+1} \equiv A_{n}^{-} F_{n-1}^{m+1} + A_{n}^{0} F_{n}^{m+1} + A_{n}^{+} F_{n+1}^{m+1} = G_{n}^{*m}, \quad n = 1(1)N_{z} - 1, \\
F_{0}^{m+1} = B_{l}(t_{m+1}), \quad F_{N_{z}}^{m+1} = B_{r}(t_{m+1}), \quad m = 0(1)N_{t} - 1, \\
F_{n}^{-s} = B_{b}(z_{n}, -t_{s}), \quad n = 0(1)N_{z}, s = 0(1)P,\n\end{cases}
$$
\n(19)

### <span id="page-5-5"></span>where

$$
A_n^- = -\frac{\varepsilon \sigma(\rho)}{h^2} + L_1 \left( -\frac{3a_{n-1}}{2h} + \frac{a_{n+1}}{2h} + p_{n-1}^{m+1} \right) - L_2 \frac{a_n}{h},
$$
  
\n
$$
A_n^0 = \frac{2\varepsilon \sigma(\rho)}{h^2} + L_1 \left( \frac{2a_{n-1}}{h} - \frac{2a_{n+1}}{h} \right) + 2L_2 p_n^{m+1},
$$
  
\n
$$
A_n^+ = -\frac{\varepsilon \sigma(\rho)}{h^2} + L_1 \left( \frac{3a_{n+1}}{2h} - \frac{a_{n-1}}{2h} + p_{n+1}^{m+1} \right) + L_2 \frac{a_n}{h},
$$
  
\n
$$
G_n^{*m} = L_1 G_{n-1}^m + 2L_2 G_n^m + L_1 G_{n+1}^m, \quad G_{n-1}^m = \frac{1}{\Delta t} F_{n-1}^m - c_{n-1}^{m+1} F_{n-1}^{m+1 - P} + g_{n-1}^{m+1},
$$
  
\n
$$
G_n^m = \frac{1}{\Delta t} F_n^m - c_n^{m+1} F_n^{m+1 - P} + g_n^{m+1}, \quad G_{n+1}^m = \frac{1}{\Delta t} F_{n+1}^m - c_{n+1}^{m+1} F_{n+1}^{m+1 - P} + g_{n+1}^{m+1}.
$$

$$
- \varepsilon \sigma(\rho) \frac{F(z_{n-1}) - 2F(z_n) + F(z_{n+1})}{h^2} + L_1 \left[ a_{n-1} F'(z_{n-1}) + p_{n-1} F(z_{n-1}) \right] + 2L_2 \left[ a_n F'(z_n) + p_n F(z_n) \right] + L_1 \left[ a_{n+1} F'(z_{n+1}) + p_{n+1} F(z_{n+1}) \right] = L_1 G_{n-1} + 2L_2 G_n + L_1 G_{n+1} + TE(h), \quad n = 1(1)N_z - 1,
$$
\n(16)

where *TE*(*h*) is local truncation error [\[37](#page-13-24)], given as

## <span id="page-5-4"></span>*Calculating ftting factor*

The fitting factor  $\sigma(\rho)$  is determined in such a way that the solution of  $(19)$  $(19)$  converges uniformly to the solution

 $h \rightarrow 0$ , we obtain

<span id="page-6-3"></span>irreducible M matrix. Hence it has positive inverse. So the existence of unique solution for  $(19)$  $(19)$  $(19)$  ensured. The

$$
\lim_{h \to 0} \frac{\sigma(\rho)}{\rho} \Big[ F_{n-1}^{m+1} - 2F_n^{m+1} + F_{n+1}^{m+1} \Big] = a_0 (L_1 + L_2) \lim_{h \to 0} \Big[ F_{n+1}^{m+1} - F_{n-1}^{m+1} \Big]. \tag{20}
$$

From singular perturbation theory [\[39](#page-13-26)] concerning to the right boundary layer, we have

<span id="page-6-0"></span>reader can refer to [\[35](#page-13-22)] for further details.  $\square$ 

<span id="page-6-1"></span>*solution*  $F_n^{m+1}$  of [\(19\)](#page-5-5), *satisfies the estimate* 

$$
F^{m+1}(z) = F_0^{m+1}(z) + (B_r(t_{m+1}) - F_0^{m+1}(1)) \exp(-a(1)(1-z)/\varepsilon) + \mathcal{O}(\varepsilon),\tag{21}
$$

where  $F_0^{m+1}(z)$  is the solution of reduced problem

**Lemma 9** (Stability result for discrete problem) The

$$
\begin{cases} \frac{F_0^{m+1}(z)-F_0^m(z)}{\Delta t} + a(z) \frac{dF_0^{m+1}(z)}{dz} + b^{m+1}(z)F_0^{m+1}(z) = -c^{m+1}(z)F_0^{m+1-P}(z) + g^{m+1}(z),\\ F_0^{-s} = B_b(z, -t_s), \quad s = 0(1)P, z \in \overline{\Omega_z}. \end{cases}
$$

Let  $0 = z_0 < z_1 < z_2 < \cdots < z_{N_z} = 1$ , such that  $z_n = nh$ ,  $n = 0(1)N_z$ . Assume *h* is sufficiently small. Then discretization of [\(21](#page-6-0)) gives

$$
F_n^{m+1} = F^{m+1}(nh) = F_0^{m+1}(nh) + (B_r(t_{m+1}) - F_0^{m+1}(1)) \exp(-a(1)(1/\varepsilon - n\rho)).
$$
\n(22)

Similarly, we have

$$
F_{n+1}^{m+1} = F_0^{m+1}((n+1)h) + (B_r(t_{m+1}) - F_0^{m+1}(1)) \exp(-a(1)(1/\varepsilon - (n+1)\rho)),
$$
  
\n
$$
F_{n-1}^{m+1} = F_0^{m+1}((n-1)h) + (B_r(t_{m+1}) - F_0^{m+1}(1)) \exp(-a(1)(1/\varepsilon - (n-1)\rho)).
$$
\n(23)

Substituting  $(22)$  $(22)$ -  $(23)$  into  $(20)$  and then simplifying, we obtain

$$
\sigma(\rho) = a(1)\rho(L_1 + L_2) \coth\left(\frac{a(1)\rho}{2}\right).
$$
 (24)

## **Stability and convergence analysis**

The uniform stability and convergence analysis for  $(19)$  $(19)$  $(19)$ are covered in this section. Firstly, we establish the existence of the unique discrete solution for the scheme ([19](#page-5-5)) by proving the discrete comparison principle.

<span id="page-6-4"></span>**Lemma 8** (Discrete Comparison Principle) *Let*   $U_n^{m+1}$  and  $F_n^{m+1}$  be two mesh functions, satisfy*ing*  $\mathscr{L}_{\varepsilon}^{\Delta t,h} U_{n}^{m+1} \leq \mathscr{L}_{\varepsilon}^{\Delta t,h} F_{n}^{m+1}, \text{ for } n = 1(1)N_{z} - 1,$  $U_0^{m+1} \le F_0^{m+1}$ , and  $U_{N_z}^{m+1} \le F_{N_z}^{m+1}$ , then  $U_n^{m+1} \le F_n^{m+1}$ *for*  $n = 0(1)N_z$ .

*Proof* The matrix associate with  $\mathscr{L}_{\varepsilon}^{\Delta t,h}$  is size of  $(N_z + 1) \times (N_z + 1)$ , where for  $n = 1$ , and  $n = N_z - 1$ , the terms demanding  $F_0^{m+1}$  and  $F_{N_z}^{m+1}$  shifted to the right side. The coefficient matrix satisfies the property of an

<span id="page-6-2"></span> $\left| F_n^{m+1} \right| \leq \pi^{-1} \left\| \mathcal{L}_{\varepsilon}^{\Delta t,h} F_n^{m+1} \right\| + \max \left\{ \left| F_0^{m+1} \right|, \left| F_{N_z}^{m+1} \right| \right\}$  $\big\}$ 

*where*  $\pi$  *is the lower bound of*  $p_n^{m+1}$ .

*Proof* L e t  $\theta = \pi^{-1} || \mathcal{L}_{\varepsilon}^{\Delta t,h} F_n^{m+1} || + \max \left\{ \left| F_0^{m+1} \right|, \left| F_{N_z}^{m+1} \right| \right\}$  $\}$ . For the barrier function  $\Theta_n^{\pm} = \theta \pm F_n^{m+1}$ , we have  $\Theta_0^{\pm} = \theta \pm F_0^{m+1} \geq 0$ ,  $\Theta_{N_z}^{\pm} = \theta \pm F_{N_z}^{m+1} \geq 0$ . Moreover for  $0 < n < N<sub>z</sub>(19)$  $0 < n < N<sub>z</sub>(19)$  can be written as

$$
\mathcal{L}_{\varepsilon}^{\Delta t, h} \Theta_{n}^{\pm} = \varepsilon \sigma(\rho) \Big[ (\theta \pm F_{n-1}^{m+1}) - 2(\theta \pm F_{n}^{m+1}) + (\theta \pm F_{n+1}^{m+1}) \Big] \n+ L_{1} \Big[ a_{n-1} (\theta \pm F_{n-1}^{m+1})' + p_{n-1}^{m+1} (\theta \pm F_{n-1}^{m+1}) \Big] \n+ 2L_{2} \Big[ a_{n} (\theta \pm F_{n}^{m+1})' + p_{n}^{m+1} (\theta \pm F_{n}^{m+1}) \Big] \n+ L_{1} \Big[ a_{n+1} (\theta \pm F_{n+1}^{m+1})' + p_{n+1}^{m+1} (\theta \pm F_{n+1}^{m+1}) \Big] \n= \Big( L_{1} p_{n-1}^{m+1} + 2L_{2} p_{n}^{m+1} + L_{1} p_{n+1}^{m+1} \Big) \theta \pm \mathcal{L}_{\varepsilon}^{\Delta t, h} F_{n}^{m+1} \n\geq 0, \quad \text{as } p_{n}^{m+1} > \pi.
$$

By applying discrete comparison principle, we obtained the required result.  $\hfill \square$ 

From the power expansion of  $\coth(q)$  and its property one can deduce that  $|q \coth(q) - 1| \leq C \frac{q^2}{q^2 + 1}$ . Thus, for  $C_1$  and  $C_2$  constants

$$
C_1 \frac{q^2}{q^2 + 1} \le q \coth(q) - 1 \le C_2 \frac{q^2}{q^2 + 1}, \text{ and } \varepsilon \frac{(h/\varepsilon)^2}{(h/\varepsilon)^2 + 1} = \frac{h^2}{h + \varepsilon}.
$$
\n(25)

Using  $(25)$  $(25)$ , we can have

$$
\left| \varepsilon \left( a(1)\rho(L_1 + L_2) \coth\left( \frac{a(1)\rho}{2} \right) - 1 \right) D^+ D^- F^{m+1}(z_n) \right| \le C \frac{h^2}{h + \varepsilon} \left\| (F^{m+1})^{(n)}(z_n) \right\|.
$$
 (26)

Then for  $n = 1(1)N_z - 1$ , using ([26\)](#page-7-1) we have

$$
| - \varepsilon(\sigma(\rho)D^{+}D^{-}F_{n}^{m+1} - (F^{m+1})^{''}(z_{n})) | = | - \varepsilon(\sigma(\rho) - 1)D^{+}D^{-}F_{n}^{m+1} - (F^{m+1})^{''}(z_{n})) |
$$
  
\n
$$
\leq \varepsilon \left| \left( a(1)\rho(L_{1} + L_{2}) \coth\left( \frac{a(1)\rho}{2} \right) -1)D^{+}D^{-}F^{m+1}(z_{n}) \right| + \varepsilon | (F^{m+1})^{''}(z_{n}) - D^{+}D^{-}F_{n}^{m+1} |
$$
  
\n
$$
\leq C \frac{h^{2}}{h + \varepsilon} \left| (F^{m+1})^{''}(z_{n}) \right| + C\varepsilon h^{2} \left| (F^{m+1})^{''}(z_{n}) \right|.
$$
  
\n(27)

<span id="page-7-2"></span>where  $\left\| (F^{m+1})^{''}(x_j) \right\| = \max_{0 \le n \le N_z} | (F^{m+1})^{''}(z_n) |$ 

 $||(F^{m+1})^{(3)}(x_j)|| = \max_{0 \le n \le N_z} |(F^{m+1})^{(3)}(z_n)|$ . The following theorem provides the spatial direction truncation error bound for the proposed scheme.

<span id="page-7-4"></span>Theorem 3 (Error in the spatial direction) Let  $a(z)$ ,  $b(z, t_{m+1})$ , andc $(z, t_{m+1})$  are sufficiently smooth *functions so that*  $F^{m+1}(z) \in C^4[\overline{\Omega_z}]$ . Then the solution  $F_n^{m+1}$  of [\(19](#page-5-5)) satisfies the estimate

$$
\left|\mathscr{L}_{\varepsilon}^{\Delta t}F^{m+1}(z_n)-\mathscr{L}_{\varepsilon}^{\Delta t,h}F_n^{m+1}\right|\leq Ch\bigg(1+\varepsilon^{-3}\exp\bigg(-\frac{\gamma(1-z_n)}{\varepsilon}\bigg)\bigg).
$$

*Proof* The local truncation error bound for ([19\)](#page-5-5) at node  $z_n$  is

From Taylor series expansion, we have

$$
|e_{n-1}'| = |(F^{m+1})'(z_{n-1}) - (F^{m+1})'_{n-1}| \leq Ch \Big\| (F^{m+1})''(x_j) \Big\|,
$$
\n(28)

<span id="page-7-3"></span>
$$
|e_{n}^{'}| = |(F^{m+1})^{'}(z_{n}) - (F^{m+1})_{n}^{'}| \leq Ch^{2} \left\| (F^{m+1})^{(3)}(x_{j}) \right\|,
$$
\n(29)

<span id="page-7-1"></span><span id="page-7-0"></span>
$$
|e_{n+1}'| = |(F^{m+1})'(z_{n+1}) - (F^{m+1})'_{n+1}| \leq Ch \Big\| (F^{m+1})''(x_j) \Big\|, \tag{30}
$$

$$
\left|\mathcal{L}_{\varepsilon}^{\Delta t}F^{m+1}(z_n) - \mathcal{L}_{\varepsilon}^{\Delta t,h}F_n^{m+1}\right| \leq \left|\left|TE(h)\right|\right| + |-\varepsilon(\sigma(\rho)D^+D^-F_n^{m+1}) - (F^{m+1})^{''}(z_n))| + |e'_{n-1}| + |e'_{n}| + |e'_{n+1}|.
$$

By bounding  $(17)$  $(17)$  and using the bounds  $(27)$  $(27)$ – $(30)$  $(30)$  we obtain Now apply Theorem [2](#page-4-6) on ([30](#page-7-3)),

$$
\left| \mathcal{L}_{\varepsilon}^{\Delta t} F^{m+1}(z_n) - \mathcal{L}_{\varepsilon}^{\Delta t, h} F_n^{m+1} \right| \leq Ch^4 (L_2 - 2L_1) \left\| (F^{m+1})^{(3)}(\varrho_n) \right\| + Ch^5(-12L_1 + 1) \left\| (F^{m+1})^{(4)}(\varrho_n) \right\| + C \frac{h^2}{h + \varepsilon} \left\| (F^{m+1})^{(n)}(z_n) \right\| + C \varepsilon h^2 \left\| (F^{m+1})^{(4)}(z_n) \right\| + L_1 Ch \left\| (F^{m+1})^{(n)}(x_j) \right\| + L_2 Ch^2 \left\| (F^{m+1})^{(n)}(x_j) \right\|.
$$
\n(31)



<span id="page-8-0"></span>(a)  $\varepsilon = 2^{-4}$ **Fig. 1** Numerical solution of Example [1](#page-11-0) with boundary layer formation,  $N = 64$ 

(b)  $\varepsilon = 2^{-12}$ 



(a)  $\varepsilon=2^{-4}$ 

(b)  $\varepsilon = 2^{-12}$ 

<span id="page-8-1"></span>**Fig. 2** Numerical solution of Example [2](#page-11-1) with boundary layer formation,  $N = 64$ 

$$
\left| \mathcal{L}_{\varepsilon}^{\Delta t} F^{m+1}(z_n) - \mathcal{L}_{\varepsilon}^{\Delta t, h} F_n^{m+1} \right| \leq Ch^4 (L_2 - 2L_1) \Big( 1 + \varepsilon^{-3} \exp(-\gamma (1 - z_n)/\varepsilon) \Big) + Ch^5 (-12L_1 + 1) \Big( 1 + \varepsilon^{-4} \exp(-\gamma (1 - z_n)/\varepsilon) \Big) + C \frac{h^2}{h + \varepsilon} \Big( 1 + \varepsilon^{-2} \exp(-\gamma (1 - z_n)/\varepsilon) \Big) + C \varepsilon h^2 \Big( 1 + \varepsilon^{-4} \exp(-\gamma (1 - z_n)/\varepsilon) \Big) + CL_1 h \Big( 1 + \varepsilon^{-2} \exp(-\gamma (1 - z_n)/\varepsilon) \Big) + CL_2 h^2 \Big( 1 + \varepsilon^{-3} \exp(-\gamma (1 - z_n)/\varepsilon) \Big) \leq Ch \Big( 1 + \varepsilon^{-3} \exp(-\gamma (1 - z_n)/\varepsilon) \Big),
$$

since  $\varepsilon^{-2} < \varepsilon^{-3}$ , and  $h^5 < h^4 < h^2 < h$ .

Now that the proof is completed.  $\hfill \square$ 

<span id="page-9-0"></span>**Lemma 10** *For a fixed mesh number*  $N_z$  *and for*  $\varepsilon \to 0$ , *we have*

<span id="page-9-1"></span>**Table [1](#page-11-0)** Maximum point wise error  $\tilde{E}^{N_Z,N_t}$ , uniform error  $\tilde{E}^{N_Z,N_t}$ ,convergence rate  $\tilde{R}^{N_Z,N_t}$ , and CPU time in second for Example 1, N  $=$  M

$\varepsilon \downarrow$	$N = 16$	$N = 32$	$N = 64$	$N = 128$	$N = 256$	$N = 512$
$2^{-0}$	4.5471e-04	2.4759e-04	1.2990e-04	6.6937e-05	3.3972e-05	1.7125e-05
$2^{-4}$	2.2371e-03	1.2671e-03	6.8474e-04	3.5796e-04	1.8316e-04	9.2701e-05
$2^{-12}$	5.9805e-03	3.6088e-03	1.9579e-03	1.0167e-03	5.1748e-04	2.5419e-04
$2^{-16}$	5.9805e-03	3.6088e-03	1.9579e-03	1.0167e-03	5.1773e-04	2.6120e-04
$2^{-20}$	5.9805e-03	3.6088e-03	1.9579e-03	1.0167e-03	5.1773e-04	2.6120e-04
$2^{-24}$	5.9805e-03	3.6088e-03	1.9579e-03	1.0167e-03	5.1773e-04	2.6120e-04
$2^{-28}$	5.9805e-03	3.6088e-03	1.9579e-03	1.0167e-03	5.1773e-04	2.6120e-04
$\tilde{F}^{N_Z,N_t}$	5.9805e-03	3.6088e-03	1.9579e-03	1.0167e-03	5.1773e-04	2.6120e-04
$\tilde{R}^{N_Z,N_t}$	7.2877e-01	8.8219e-01	9.4545e-01	9.7359e-01	9.8703e-01	
CPU time (s)	0.0536	0.0659	0.0897	0.1120	0.4506	2.7989

<span id="page-9-2"></span>**Table 2** Maximum point wise error  $\tilde{E}^{N_Z,N_t}$ , uniform error  $\tilde{E}^{N_Z,N_t}$ , convergence rate  $\tilde{R}^{N_Z,N_t}$ , and CPU time in second for Example [2,](#page-11-1) N  $=$  M



$\varepsilon \downarrow$	$N = 16$	$N = 32$ $M = 40$	$N = 64$ $M = 80$	$N = 128$ $M = 160$	$N = 256$ $M = 320$	$N = 512$ $M = 640$
	$M = 20$					
Proposed Method						
$2^{-0}$	3.7207e-04	2.0193e-04	1.0499e-04	5.3852e-05	2.7270e-05	1.3722e-05
$2^{-4}$	1.8438e-03	1.0392e-03	5.5544e-04	2.8862e-04	1.4715e-04	7.4318e-05
$2^{-12}$	5.6785e-03	3.4308e-03	1.8570e-03	9.6293e-04	4.8950e-04	2.3987e-04
$2^{-16}$	5.6785e-03	3.4308e-03	1.8570e-03	9.6293e-04	4.8975e-04	2.4688e-04
$2^{-20}$	5.6785e-03	3.4308e-03	1.8570e-03	9.6293e-04	4.8975e-04	2.4688e-04
$2^{-24}$	5.6785e-03	3.4308e-03	1.8570e-03	9.6293e-04	4.8975e-04	2.4688e-04
$2^{-28}$	5.6785e-03	3.4308e-03	1.8570e-03	9.6293e-04	4.8975e-04	2.4688e-04
$2^{-32}$	5.6785e-03	3.4308e-03	1.8570e-03	9.6293e-04	4.8975e-04	2.4688e-04
Results in [40]						
$2^{-0}$	1.95e-04	1.04e-04	5.35e-05	2.72e-05	1.37e-05	6.88e-06
$2^{-4}$	2.00e-03	6.71e-04	3.42e-04	2.27e-04	$1.21e-04$	5.39e-05
$2^{-12}$	1.96e-02	1.03e-02	5.24e-03	2.64e-03	1.32e-03	6.60e-04
$2^{-16}$	1.97e-02	1.04e-02	5.27e-03	2.64e-03	1.32e-03	6.66e-04
$2^{-20}$	1.97e-02	1.04e-02	5.28e-03	2.64e-03	1.32e-03	6.59e-04
$2^{-24}$	1.97e-02	1.04e-02	5.28e-03	2.64e-03	1.32e-03	6.60e-04
$2^{-28}$	1.97e-02	1.04e-02	5.28e-03	2.64e-03	1.32e-03	6.60e-04
$2^{-32}$	1.97e-02	1.04e-02	5.28e-03	2.64e-03	1.32e-03	6.60e-04

<span id="page-10-0"></span>**Table 3** Comparison of maximum point wise errors for Example [1](#page-11-0)

<span id="page-10-1"></span>**Table 4** Comparison of maximum point wise errors for Example [2](#page-11-1)

$\varepsilon \downarrow$	$N = 16$	$N = 32$	$N = 64$	$N = 128$	$N = 256$	$N = 512$
Proposed Method						
$2^{-0}$	4.7516e-04	2.3325e-04	1.1624e-04	5.8105e-05	2.9052e-05	1.4526e-05
$2^{-4}$	2.7847e-03	1.0798e-03	6.4791e-04	3.5894e-04	1.8853e-04	9.6561e-05
$2^{-8}$	7.2234e-03	3.6109e-03	1.6076e-03	5.5890e-04	1.6922e-04	8.6042e-05
$2^{-12}$	7.2237e-03	3.6438e-03	1.8325e-03	9.1984e-04	4.6076e-04	2.2852e-04
$2^{-16}$	7.2237e-03	3.6438e-03	1.8325e-03	9.1984e-04	4.6078e-04	2.3061e-04
$2^{-20}$	7.2237e-03	3.6438e-03	1.8325e-03	9.1984e-04	4.6078e-04	2.3061e-04
$2^{-24}$	7.2237e-03	3.6438e-03	1.8325e-03	9.1984e-04	4.6078e-04	2.3061e-04
$2^{-28}$	7.2237e-03	3.6438e-03	1.8325e-03	9.1984e-04	4.6078e-04	2.3061e-04
Results in [41]						
$2^{-0}$	2.3950e-02	1.7664e-02	1.1228e-02	6.4886e-03	3.5334e-03	1.8536e-03
$2^{-4}$	4.8048e-02	2.7869e-02	1.4847e-02	7.6292e-03	3.8619e-03	1.9422e-03
$2^{-8}$	4.9006e-02	2.8622e-02	1.5142e-02	7.7170e-03	3.8852e-03	1.9482e-03
$2^{-12}$	4.9006e-02	2.8622e-02	1.5141e-02	7.7173e-03	3.8858e-03	1.9484e-03
$2^{-16}$	4.9006e-02	2.8622e-02	1.5141e-02	7.7173e-03	3.8858e-03	1.9484e-03
$2^{-20}$	4.9006e-02	2.8622e-02	1.5141e-02	7.7173e-03	3.8858e-03	1.9484e-03
$2^{-24}$	4.9006e-02	2.8622e-02	1.5141e-02	7.7173e-03	3.8858e-03	1.9484e-03
$2^{-28}$	4.9006e-02	2.8622e-02	1.5141e-02	7.7173e-03	3.8858e-03	1.9484e-03



<span id="page-11-3"></span>**Fig. 3** Log-log plot of maximum point wise error [3](#page-11-3) **a** for Example [1](#page-11-0) and [3](#page-11-3) **B** for Example [2](#page-11-1)

$$
\lim_{\varepsilon \to 0} \max_{1 \le n \le N_{\varepsilon}-1} \varepsilon^{-\lambda} \exp\left(-\frac{\gamma(1-z_n)}{\varepsilon}\right) = 0, \lambda \text{ is positive integer.}
$$

*Proof* Refer [\[14\]](#page-13-13). □

<span id="page-11-2"></span>**Theorem 4** *Let*  $F^{m+1}(z_n)$  *and*  $F_n^{m+1}$  *be the solutions of* ([5\)](#page-3-1) *and* ([19\)](#page-5-5) *respectively*, *then we have uniform error estimate*

$$
\left\|F^{m+1}(z_n) - F_n^{m+1}\right\| = \sup_{0 < \varepsilon < 1} |F^{m+1}(z_n) - F_n^{m+1}| \leq Ch.
$$

$$
\begin{cases}\n\frac{\partial f}{\partial t} - \varepsilon \frac{\partial^2 f}{\partial z^2} + (2 - z^2) \frac{\partial f}{\partial z} + (z + 1)(t + 1)f(z, t) \\
= -f(z, t - 1) + 10t^2 \exp(-t)z(1 - z), (z, t) \in (0, 1) \times (0, 2], \\
f(0, 1) = 0, f(1, t) = 0, t \in (0, 2], \\
f(z, t) = 0, (z, t) \in [0, 1] \times [-1, 0],\n\end{cases}
$$

whose exact solution is not known.

<span id="page-11-1"></span>*Example 2* Now consider the following problem [\[41\]](#page-13-28):

$$
\begin{cases}\n\frac{\partial f}{\partial t} - \varepsilon \frac{\partial^2 f}{\partial z^2} + \frac{(5-z^2)}{3} \frac{\partial f}{\partial z} + tf(z, t) \\
= -f(z, t-1) + t^3 z(1-z) \sin(\pi z), (z, t) \in (0, 1) \times (0, 2], \\
f(0, 1) = 0, f(1, t) = 0, t \in (0, 2], \\
f(z, t) = 0, (z, t) \in [0, 1] \times [-1, 0],\n\end{cases}
$$

whose exact solution is not known.

Figures [1](#page-8-0) and [2](#page-8-1) show the physical behavior surface graph of the numerical solutions to Examples [1](#page-11-0) and [2](#page-11-1), respectively. From these graphs, we observe that the solution exhibits a boundary layer near  $x = 1$  for

*Proof* Use the result in Lemma [10](#page-9-0) to Theorem [3,](#page-7-4) then applying Lemma [8](#page-6-4) gives the required result.  $\Box$ 

The uniform error bound of the scheme in the maximum norm is provided by the following theorem.

**Theorem 5** Let  $f(z_n, t_m)$  and  $F_n^m$  are to be the solution *of the continuous problem* ([1\)](#page-2-0) *and the discrete problem*  ([19\)](#page-5-5) respectively. Then we have

$$
||f(z_n, t_m) - F_n^m|| \le C(\Delta t + h), \quad n = 0(1)N_z, m = 0(1)N_t.
$$

*Proof* The proof follows from the combination of Lemma 7 and Theorem 4 Lemma [7](#page-4-3) and Theorem [4.](#page-11-2)

## **Numerical examples, results, and discussion**

We consider two problems of singularly perturbed parabolic diferential equations with large delays in order to illustrate the applicability of the method.

<span id="page-11-0"></span>*Example 1* Consider the following problem [[40](#page-13-27)]:

different values of  $\varepsilon$ . The graphs also display the effect of the perturbation parameter  $\varepsilon$ . That is, as the values of  $\varepsilon$  decrease, the width of the boundary layer decreases. The log-log plot in Figs. [3](#page-11-3) also approves the  $\varepsilon$ -uniform convergence of the proposed method. It is evident from these fgures that as mesh numbers increases, the maximum point wise error decreases monotonically.

As the exact solutions to the problems are not known, we use the double mesh technique [[42](#page-13-29)] to compute the maximum pointwise error  $\tilde{E}_{\varepsilon}^{N_z, N_t}$  and order of convergence  $\tilde{R}_{\varepsilon}^{N_z,N_t}$  as follows

$$
\tilde{E}_{\varepsilon}^{N_z, N_t} = \max_{\substack{0 \le j \le N_z \\ 0 \le i \le N_t}} \left| F^{N_z, N_t}(z_j, t_i) - F^{2N_z, 2N_t}(z_j, t_i) \right|, \n0 \le i \le N_t \n\tilde{R}_{\varepsilon}^{N_z, N_t} = \log_2 \tilde{E}_{\varepsilon}^{N_z, N_t} - \log_2 \tilde{E}_{\varepsilon}^{2N_z, 2N_t},
$$
\n(32)

where  $F^{N_z, N_t}(z_i, t_i)$  and  $F^{2N_z, 2N_t}(z_i, t_i)$  are the numerical approximation of the exact solution  $f(z_j, t_i)$  on the mesh  $\Omega^N$  and  $\Omega^{2N}$  respectively.  $\Omega^{2N}$  is obtained by doubling the mesh  $\Omega^N$  such that the mid points  $z_{i+1/2} = (z_{i+1} + z_i)/2$ and  $t_{i+1/2} = (t_{i+1} + t_i)/2$  are included in to the mesh points. From ([32\)](#page-12-0), we compute  $\varepsilon$ -uniform maximum error  $\tilde{E}^{N_z,N_t}$  and the corresponding  $\varepsilon$ -uniform convergence rate  $\tilde{R}^{N_{z},N_{t}}$  as

$$
\tilde{E}^{N_z,N_t} = \max_{\varepsilon} \tilde{E}_{\varepsilon}^{N_z,N_t},
$$
\n
$$
\tilde{R}^{N_z,N_t} = \log_2 \tilde{E}^{N_z,N_t} - \log_2 \tilde{E}^{2N_z,2N_t}.
$$
\n(33)

The numerical results presented in Tables [1](#page-9-1) and [2](#page-9-2) show the maximum point-wise error, ε-uniform maximum error, rate of convergence, and CPU time in seconds for the proposed method for Examples [1](#page-11-0) and [2](#page-11-1), respectively. From these tables, we confrm that the suggested method is linear-order ε-uniformly convergent, according to the error analyses carried out in this work. Furthermore, we see that as the mesh numbers increase, the maximum point-wise error decreases, and as the values of  $\varepsilon$  decrease, a stable and bounded maximum error is established. Thus, the proposed scheme is  $\varepsilon$ -uniformly convergent.

Tables [3](#page-10-0) and [4](#page-10-1) show the comparison of the maximum point wise error of the methods that existed in the literature [\[40](#page-13-27), [41\]](#page-13-28) for Examples 1 and 2, respectively. In Table [3,](#page-10-0) as the  $\varepsilon$  gets smaller, we observe that the proposed method holds a more accurate  $\varepsilon$  uniform convergence than the method in [\[40](#page-13-27)]. Similarly, in Table [4](#page-10-1), the comparison shows that the results of the proposed scheme are more accurate  $\varepsilon$  uniform convergence than the method in  $[41]$  $[41]$ .

## **Conclusion**

In this work, a non-classical numerical method is developed to solve a class of singularly perturbed delay parabolic convection-difusion problems with Dirichlet boundary conditions. The solutions of these problems display boundary layer at the right side of spatial domain as  $\varepsilon \to 0$ . A delay term is handled by constructing a mesh in such a way that the delay argument coincides with a mesh point. The method is based on exponential spline on uniform mesh with exponential ftting factor. It is shown that the method is uniformly convergent independent of mesh parameters and perturbation parameter and provides uniform first-order convergence. The proposed method has the advantage of being applicable to personal computers with very low CPU computing time for the required number of mesh points. For the proposed method the evaluation of exponential functions is computationally intensive, particularly if high precision is needed. This difficulty can add to the overall CPU time required for solving the differential equations. The theoretical result is validated numerically by two test examples that are presented. The performance of the method was compared with some existing literatures and gave more accurate result.

#### <span id="page-12-0"></span>**Abbreviations**



#### **Acknowledgements**

The authors are grateful to the anonymous referees and editors for their constructive comments.

#### **Author contributions**

ZIH designed, analysis, drafted the work, coding MATLAB and numerical experimentation. GFD designed, analysis, drafted the work. Both authors read and approved the fnal manuscript.

## **Funding**

There is no any fund support in this research work.

#### **Availability of data and materials**

There is no additional data used for this research work.

#### **Declarations**

**Ethics approval and consent to participate** Not applicable.

#### **Consent for publication**

Not applicable.

#### **Competing interests**

The authors declare that the research was conducted in the absence of any commercial or fnancial relationships that could be construed as a potential Confict of interest

Received: 21 January 2024 Accepted: 14 November 2024

#### **References**

- <span id="page-13-0"></span>Glizer V. Asymptotic solution of a boundary-value problem for linear singularly-perturbed functional diferential equations arising in optimal control theory. J Optim Theory Appl. 2000;106:309–35.
- <span id="page-13-1"></span>2. Vulanovic R, Farrell PA, Lin P. Numerical solution of nonlinear singular perturbation problems modeling chemical reactions. Applications of advanced computational methods for boundary and interior layers. 1993;192–213.
- <span id="page-13-2"></span>3. Salpeter EE, Salpeter SR. Mathematical model for the epidemiology of tuberculosis, with estimates of the reproductive number and infectiondelay function. Am J Epidemiol. 1998;147(4):398–406.
- <span id="page-13-3"></span>Mallet-Paret J, Nussbaum RD. A differential-delay equation arising in optics and physiology. SIAM J Math Anal. 1989;20(2):249–92.
- <span id="page-13-4"></span>5. Campbell SA, Edwards R, Driessche P. Delayed coupling between two neural network loops. SIAM J Appl Math. 2004;65(1):316–35.
- <span id="page-13-5"></span>6. Roos H-G. Robust numerical methods for singularly perturbed diferential equations. Berlin: Springer; 2008.
- <span id="page-13-6"></span>7. Kaushik A, Sharma M. Convergence analysis of weighted diference approximations on piecewise uniform grids to a class of singularly perturbed functional diferential equations. J Optim Theory Appl. 2012;155(1):252–72.
- <span id="page-13-7"></span>8. Das A, Natesan S. Uniformly convergent hybrid numerical scheme for singularly perturbed delay parabolic convection-difusion problems on shishkin mesh. Appl Math Comput. 2015;271:168–86.
- <span id="page-13-8"></span>9. Gowrisankar S, Natesan S.  $\varepsilon$ -uniformly convergent numerical scheme for singularly perturbed delay parabolic partial diferential equations. Int J Comput Math. 2017;94(5):902–21.
- <span id="page-13-9"></span>10. Woldaregay MM, Aniley WT, Duressa GF. Novel numerical scheme for singularly perturbed time delay convection-difusion equation. Adv Math Phys. 2021;2021:1–13.
- <span id="page-13-10"></span>11. Negero NT, Duressa GF. A method of line with improved accuracy for singularly perturbed parabolic convection-difusion problems with large temporal lag. Results Appl Math. 2021;11: 100174.
- <span id="page-13-11"></span>12. Negero NT, Duressa GF. An efficient numerical approach for singularly perturbed parabolic convection-difusion problems with large time-lag. J Math Model. 2022;10(2):173.
- <span id="page-13-12"></span>13. Negero NT, Duressa GF. An exponentially ftted spline method for singularly perturbed parabolic convection-difusion problems with large time delay. Tamkang J Math. 2023;54(4):313–38.
- <span id="page-13-13"></span>14. Hassen ZI, Duressa GF. New approach of convergent numerical method for singularly perturbed delay parabolic convection-difusion problems. Res Math. 2023;10(1):2225267.
- <span id="page-13-14"></span>15. Tesfaye SK, Woldaregay MM, Dinka TG, Duressa GF. Fitted computational method for solving singularly perturbed small time lag problem. BMC Res Notes. 2022;15(1):1–10.
- <span id="page-13-15"></span>16. Tesfaye SK, Duressa GF, Woldaregay MM, Dinka T. Fitted computational method for singularly perturbed convection-difusion equation with time delay. Front Appl Math Stat. 2023;9:1244490.
- <span id="page-13-16"></span>17. Hassen ZI, Duressa GF. Parameter-uniformly convergent numerical scheme for singularly perturbed delay parabolic diferential equation via extended B-spline collocation. Front Appl Math Stat. 2023;9:1255672.
- <span id="page-13-17"></span>18. Kumar A, Gowrisankar S. Efficient numerical methods on modified graded mesh for singularly perturbed parabolic problem with time delay. Iran J Numerical Anal Opt. 2024;14(1):77–106.
- <span id="page-13-18"></span>19. Kumar D, Kumari P. A parameter-uniform scheme for singularly perturbed partial diferential equations with a time lag. Numerical Methods Partial Diff Equ. 2020;36(4):868-86.
- 20. Govindarao L, Sahu SR, Mohapatra J. Uniformly convergent numerical method for singularly perturbed time delay parabolic problem with two small parameters. Iran J Sci Technol Trans A Sci. 2019;43:2373–83.
- 21. Rai P, Yadav S. Robust numerical schemes for singularly perturbed delay parabolic convection-diffusion problems with degenerate coefficient. Int J Comput Math. 2021;98(1):195–221.
- 22. Yadav S, Rai P. A higher order scheme for singularly perturbed delay parabolic turning point problem. Eng Comput. 2021;38(2):819–51.
- 23. Kumar K, Podila PC, Das P, Ramos H. A graded mesh refnement approach for boundary layer originated singularly perturbed time-delayed parabolic convection difusion problems. Math Methods Appl Sci. 2021;44(16):12332–50.
- 24. Woldaregay MM, Hunde TW, Mishra VN. Fitted exact diference method for solutions of a singularly perturbed time delay parabolic pde. Partial Dif Equ Appl Math. 2023;8: 100556.
- 25. Das A, Natesan S. Second-order uniformly convergent numerical method for singularly perturbed delay parabolic partial diferential equations. Int J Comput Math. 2018;95(3):490–510.
- 26. Hassen ZI, Duressa GF. Parameter uniform hybrid numerical method for time-dependent singularly perturbed parabolic diferential equations with large delay. Appl Math Sci Eng. 2024;32(1):2328254.
- 27. Hassen ZI, Duressa GF. Nonstandard hybrid numerical scheme for singularly perturbed parabolic diferential equations with large delay. Partial Diff Equ Appl Math. 2024;10: 100722.
- 28. Hailu WS, Duressa GF. Uniformly convergent numerical scheme for solving singularly perturbed parabolic convection-difusion equations with integral boundary condition. Dif Equ Dyn Syst. 2023;1–27.
- 29. Amiraliyev GM, Cimen E, Amirali I, Cakir M. High-order fnite diference technique for delay pseudo-parabolic equations. J Comput Appl Math. 2017;321:1–7.
- 30. Gunes B, Duru H. A computational method for the singularly perturbed delay pseudo-parabolic diferential equations on adaptive mesh. Int J Comput Math. 2023;100(8):1667–82.
- 31. Babu G, Prithvi M, Sharma KK, Ramesh V. A uniformly convergent numerical algorithm on harmonic ( $H(\ell)$ ) mesh for parabolic singularly perturbed convection-difusion problems with boundary layer. Dif Equ Dyn Syst. 2024;32(2):551–64.
- <span id="page-13-19"></span>32. Howlader J, Mishra P, Sharma KK. An orthogonal spline collocation method for singularly perturbed parabolic reaction-difusion problems with time delay. J Appl Math Comput. 2024;70(2):1069–101.
- <span id="page-13-20"></span>33. Protter MH, Weinberger HF. Maximum Principles in Diferential Equations. Erscheinungsort nicht ermittelbar: Springer; 2012.
- <span id="page-13-21"></span>34. Mbroh NA, Noutchie SCO, Massoukou RYM. A robust method of lines solution for singularly perturbed delay parabolic problem. Alex Eng J. 2020;59(4):2543–54.
- <span id="page-13-22"></span>35. Kellogg RB, Tsan A. Analysis of some diference approximations for a singular perturbation problem without turning points. Math Comput. 1978;32(144):1025–39.
- <span id="page-13-23"></span>36. Zahra WK. Finite-diference technique based on exponential splines for the solution of obstacle problems. Int J Comput Math. 2011;88(14):3046–60.
- <span id="page-13-24"></span>37. Adivi Sri Venkata RK, Palli MMK. A numerical approach for solving singularly perturbed convection delay problems via exponentially ftted spline method. Calcolo. 2017;54:943–61.
- <span id="page-13-25"></span>38. Ranjan R, Prasad HS. A novel approach for the numerical approximation to the solution of singularly perturbed diferential-diference equations with small shifts. J Appl Math Comput. 2021;65(1–2):403–27.
- <span id="page-13-26"></span>39. O'Malley RE. Introduction to Singular Perturbations. Applied mathematics and mechanics. New York: Academic Press; 1974.
- <span id="page-13-27"></span>40. Kumar D. A parameter-uniform scheme for the parabolic singularly perturbed problem with a delay in time. Numerical Methods Partial Diff Equ. 2021;37(1):626–42.
- <span id="page-13-28"></span>41. Babu G, Bansal K. A high order robust numerical scheme for singularly perturbed delay parabolic convection difusion problems. J Appl Math Comput. 2021;1–27.
- <span id="page-13-29"></span>42. Johnson C. Uniform Numerical Methods for Problems with Initial and Boundary Layers (EP Doolan, JJH Miller and WHA Schilders). Society for Industrial and Applied Mathematics (1983)

## **Publisher's Note**

Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.